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Designing Dimeric Lanthanide(III)-Containing Ionic liquids

Éadaoin McCourt,^[a] Kane Esien,^[b] Li Zhenyu,^[a] Solveig Felton,^{*[b]} and Peter Nockemann^{*[a]}

Abstract: Herein, we report on the preparation of liquid dimeric lanthanide(III)-containing compounds. Starting from the design of dimeric solids, we demonstrate that by tuning of anion and cation structures we can lower the melting points below room temperature, whilst maintaining the dimeric structure. Magnetic measurements could establish the spin-spin interactions of the neighboring lanthanide(III) ions in the liquid state at low temperatures, and matched the interactions of the analogous crystalline solid compounds.

Lanthanide-containing soft materials are a fascinating class of materials with intrinsic luminescent^[1] and magnetic properties.^[2] The potential of lanthanides to generate significant magnetic anisotropy has led to di- or polynuclear lanthanide (III) complexes receiving great attention.^[3-6] Polymetallic assemblies can exhibit cooperative metalmetal interactions affecting their magnetic, electronic and photophysical properties.[7]

lonic liquids (ILs) by definition are salts which have melting points below 100 °C or commonly at room temperature (RTILs).^[8] Properties such as low-volatility, high conductivity, extensive anion and cation tunability, wide liquid range and high thermal stability are found for these compounds.^[9–12] The first ionic liquid reported to show a strong response to a magnet was 1-butyl-3-methylimidazolium tetrachloroferrate, [C₄Mim][FeCl₄].^[13,14]

Since then, many magneto-active metal-containing ionic liquids have been reported,^[15] in which the magnetic functionality is incorporated into the anion structure, e.g. Mn^{II}Cl₄² Co^{II}(NCS)₄^{2. [16-18]} Magnetic measurements carried out on [C₂Mim][FeCl₄] have also shown long range antiferromagnetic ordering at low temperatures.[19]

We have previously reported on the first liquid complexes containing lanthanide anions with thiocyanate ligands.^[20] The magnetic and luminescent properties of [C₆Mim]_{5-x}[Dy(SCN)_{8-x}(H₂O)_x] where subsequently explored by Mallick et al.^[2] Single-ion magnetic behavior at low temperature has been reported for the Dy-containing

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IL $[Dy(R-Im-CH_2COO)_3(H_2O))_2][(PF_6)_3 \cdot 2H_2O].^{[21]}$ In this case the intrinsic magnetic functionality was incorporated into the cation structure. Recently, we investigated the physical and magnetic properties of the series [P_{666 14}][RECl₆] (RE = Rare Earth).^[22] Further low-temperature magnetic measurements on [P_{666 14}][RECl₆] ILs were subsequently carried out by Monson et al.,[23] reporting unusual behavior at temperatures below 50 K, attributed to the trapping of intermediate structures during glass formation.

In this work, the magnetic and structural properties of novel crystalline, glassy and liquid lanthanide(III) containing ionic compounds, $[C_4Mim]_2[Ln_2(CH_3COO)_8]$, $[P_{666\ 14}]_2[Ln_2(CH_3COO)_8]$ and $[P_{666 \ 14}]_2[Ln_2(C_7H_{15}COO)_8]$ are reported (Ln = lanthanide series). We demonstrate the ability to fine-tune both the anion and cation structures to realize room-temperature lanthanide-containing ionic liquids shown in Figure 1.



Figure 1. Room-temperature lanthanide-containing ionic liquids, $[P_{666 \ 14}]_2[Ln_2(C_7H_{15}COO)_8]$ Ln = lanthanide.

This paper will focus on the dysprosium (III) containing crystalline, glassy and liquid complexes. Single-crystal X-ray diffraction provided direct evidence of the formation of a dimeric Dy...Dy complex in the solid state [C₄Mim]₂[Dy₂(CH₃COO)₈] Figure 2.



Figure 2. Crystal structure of the dimeric [C₄Mim]₂[Dy₂(CH₃COO)₈].

This ultimately led to the questions (i) "Can a dimeric structure be obtained in the liquid state?" and, if so, (ii) "How can the presence of a dimer be proven in a liquid?".

These questions are the underlying foundation of this paper and we endeavor to answer them through a combination of techniques discussed herein. In particular, SQUID (Superconducting Quantum Interference Device) magnetometry enabled us to explore the magnetic properties of these three systems bestowing insights into the magnetic behavior over a wide temperature range. Although multiple-technique approaches, including HEXS (High Energy X-ray Scattering) and EXAFS (Extended X-ray Absorption Fine Structure) to probe the speciation of metals in the liquid state have been previously reported, the methods may provide limited information and are not easily accessible.^[24] Therefore, we looked for the presence of a magnetic interaction between the neighboring dysprosium ions. Comparison between the magnetic properties of the three samples, if such an interaction is present, will allow us to infer structure in the liquid and glassy states.

The novel lanthanide(III)-containing ionic compounds, $[C_4Mim]_2[Ln_2(CH_3COO)_8]$, $[P_{666\ 14}]_2[Ln_2(CH_3COO)_8]$ and $[P_{666\ 14}]_2[Ln_2(C_7H_{15}COO)_8]$ were synthesized *via* addition reactions, whereby, equimolar ratios of lanthanide(III) salt (acetate or octanoate) were reacted with the ionic liquids 1-butyl-3-methylimidazolium acetate, trihexyltetra(decyl)phosphonium acetate and trihexytetra(decyl)phosphonium octanoate. By careful selection of the cation and anion components the melting points could be tuned forming the crystalline, glass and liquid lanthanide(III)-containing compounds shown in Figure 4.

By a method of slow crystallization in a Reacti-Therm, starting at 120 °C from the molten state with a cooling rate of -5 °C (m.p. = 118 °C) to 50 °C, single crystals of [C₄Mim]₂[Dy₂(CH₃COO)₈] formed. The structure presented in Figure 2, consists of a [Dy₂(CH₃COO)₈]²⁻ anion and two 1-butyl-3-methylimidazolium, [C₄Mim]⁺ counterions. Powder XRD was used to confirm that what was obtained in the single crystal was indicative of the bulk sample (Figure S1). The Dy^{III} ions are linked by two carboxylate ligands in a bridging bidentate $(\eta^1:\eta^1:\mu_2)$ fashion and two in a bridging chelating $(\eta^1:\eta^2:\mu_2)$ mode. Each Dy^{III} ion is further surrounded by two chelating carboxylates giving a total coordination number of nine. As previously discussed by Ouchi et al.^[6] in the bridging chelating carboxylate groups, one of the oxygen atoms is bound in a μ_2 -bridging manner to both dysprosium atoms and the second is bound directly to one dysprosium atom as shown in Figure 2. The coordination polyhedron can be described as a distorted tricapped trigonal prism with Dy - O (carboxylate) bond lengths ranging from 2.329 - 2.509 Å and the Dy---Dy distance within the dimer is 3.8587(6) Å (for Dy1) and 3.8872(6) Å (for Dy2), respectively.

As diffraction experiments are limited to the crystalline sample other spectroscopic techniques such as IR and UV-Vis were

investigated, however, whilst affording some insight into the structure, the standard spectroscopic techniques did not provide hard evidence for the question *(ii)* (*vide supra*) regarding the dimeric structure in the liquid state. The carboxylate ion, RCO_2^- , has several possible coordination modes to a metal *e.g.* monodentate, bidentate chelating, bidentate bridging. Analysis of the differences (Δv) in the asymmetric $v_a(COO^-)$ and symmetric $v_s(COO^-)$ wavenumbers, can be used to gain insight into the coordination modes, [$\Delta v = v_a(COO^-) - v_s(COO^-)$].^[25] From the spectra and values presented in Figure 3 and Table 1, splitting of $v_a(COO^-)$ is observed in the crystalline, glassy and liquid samples.



Figure 3. Infrared spectra of [C₄Mim]₂[Dy₂(CH₃COO)₈], [P_{666 14}]₂[Dy₂(CH₃COO)₈] and [P_{666 14}]₂[Dy₂(C₇H₁₅COO)₈].

This splitting leads to two separation values indicative of different bonding modes present within the sample. A mix of bidentate bridging and chelating may be deduced for the crystalline, glassy and liquid samples. UV-Vis spectroscopy of $[P_{666\ 14}]_2[Dy_2(C_7H_{15}COO)_8]$ (Figure S12) showed an absorbance spectrum characteristic of Dy^{III} , whilst thermal gravimetrical analysis (TGA) and differential scanning calorimetry (DSC) provided information about decomposition temperatures and phase transitions (Table 2), respectively.

Table 1: Values of $v_a(COO^-)$, $v_s(COO^-)$ and $\Delta_v(COO^-)$.

Compound	v _a (COO⁻) (cm⁻¹)	v s(COO ⁻) (cm ⁻¹)	∆v ^[a]
Na⁺ CH₃COO⁻	1578	1414	164 ^[25]
[C ₄ Mim] ₂ [Dy ₂ (CH ₃ COO) ₈]	1598/ 1564	1431	167/ 133
[P _{666 14}] ₂ [Dy ₂ (CH ₃ COO) ₈]	1578/ 1564	1435	143/ 129
[P _{666 14}] ₂ [Dy ₂ (C ₇ H ₁₅ COO) ₈]	1581/ 1558	1435	146/ 123

 $[a] \Delta_v = v_a(COO^-) - v_s(COO^-)$



Figure 4. Variations in anion and cation structure upon changing from crystalline solid to glass to liquid, with $a \rightarrow b$ changing cation from [C₄Mim]⁺ to [P_{666 14}]⁺ and $b \rightarrow c$ changing coordinating anion from [OAc]⁻ to [C₇H₁₅COO]⁻ (Ln = lanthanide). Illustration of the ordered state in the solid (left), with Coulombic attraction dominating and the transition via a glass state into a liquid (right) with predominantly van der Waals interactions.

Table 2.	Thermal	properties	of the	crystalline,	glassy a	and liquid	samples.
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Compound	T _m ^a (°C)	Т _g ^b (°С)	T _{dec} ^c (°C)
[C ₄ Mim] ₂ [Dy ₂ (CH ₃ COO) ₈]	118	-	292
[P _{666 14}] ₂ [Dy ₂ (CH ₃ COO) ₈]	-	-41.3	327
[P _{666 14}]2[Dy ₂ (C ₇ H ₁₅ COO)8]	-	-80.4	354

^a Melting point (T_m), ^b Glass transition (T_g), ^c Onset of decomposition (T_{dec}).

These techniques combined give an overall insight into the physical properties of the studied systems. On their own, they are inadequate in confirming the hypothesis that the dimeric Dy … Dysub-units established in the solid state are evident in the glass and liquid state. To gain a more in-depth understanding of the glass and liquid properties, the intrinsic magnetic properties of all three samples were investigated.

A superconducting quantum interference device (SQUID) magnetometer was employed to monitor the magnetic moment of each compound over a wide temperature range; it was hoped that an antiferromagnetic or ferromagnetic signature would be observed in the magnetic susceptibility of each respective compound, indicative of magnetic dimers beginning to order. The susceptibility and inverse susceptibility as a function of temperature for the crystalline solid, liquid and glass compounds, are available in the supplementary material (Figure S13) along with the results from Curie-Weiss fits to these data.

The susceptibility multiplied by the temperature as a function of temperature (χ T), for the crystalline solid, liquid and glass compounds is presented in Figure 5. Here clear deviation from non-interacting behavior below 100 K for each compound is seen. The χ T value that

is predicted for two non-interacting Dy^{III} ions (⁶H_{15/2} with g = 4/3^[26]) per formula unit is shown in Figure 5 by the dashed line. Considering the constant high temperature regions, it was found that the solid crystalline sample and the liquid sample yielded effective magnetic moments per Dy^{III} ion of 10.7 μ_B and 10.4 μ_B , respectively. These values agree very well with the predicted value of 10.6 μ_B , corresponding to two Dy^{III} ions per formula unit in the ground state ⁶H_{15/2} with g = 4/3.^[26] Regarding experimentally measured magnetic moments of Dy^{III} ions, a range or 10.2 – 10.6 μ_B per Dy^{III} ion is observed in the literature^[27–29], with our previous work investigating monomeric lanthanide-based ionic liquids yielding a value of 10.2 μ_B per Dy^{III} ion.^[22]

We can conclude from Figure 5 that above 100 K, we see no evidence of any interactions between Dy^{III} ions. Below 100 K the measured xT values begin to deviate from that predicted for non-interacting magnetic moments. For DyIII ions, one cannot simply assume that a decrease in χT values directly corresponds to an antiferromagnetic coupling; the possibility that this decrease is due to a combination of dysprosium's magnetic anisotropy, progressive depopulation of the Stark levels - arising from crystal field splitting - and the exchange interaction^[30] must be entertained. A full numerical treatment isn't viable due to the large anisotropy of the Dy^{III} ion. However, it is possible to make a rough quantitative model of the system. This method was first introduced by McPherson et al.[31] and has been widely used to describe the magnetic behavior of lanthanide containing chemical compounds.^[5,32-35] The crux of the model is that the magnetic energy levels (m_J) are considered LS coupled energy levels which are zero-field-split in an axial field geometry $(\hat{H} = \Delta \hat{J}_z^2)$ magnetic susceptibility of each Dy^{III} ion is given by Equation 1:

$$\chi_{Dy} = \frac{Ng^2\beta^2}{4KT} \frac{A}{B}$$
(1)

$$A = 225e \frac{-225\Delta}{4KT} + 169e \frac{-169\Delta}{4KT} + 121e \frac{-121\Delta}{4KT} + 81e \frac{-81\Delta}{4KT} + 49e \frac{-49\Delta}{4KT} + 25e \frac{-25\Delta}{4KT} + 9e \frac{-9\Delta}{4KT} + e \frac{-\Delta}{4KT}$$

 $B = e \frac{-225\Delta}{4KT} + e \frac{-169\Delta}{4KT} + e \frac{-121\Delta}{4KT} + e \frac{-81\Delta}{4KT} + e \frac{-49\Delta}{4KT} + e \frac{-25\Delta}{4KT} + e \frac{-9\Delta}{4KT} + e \frac{-\Delta}{4KT}$

The interactions between magnetic moments are simulated by incorporating the Weiss parameter according to Equation 2:

$$\chi_M = \frac{T}{T - \theta} 2 \chi_{Dy} \quad (2)$$

Through fitting the data, it is possible to extract numerical values for the zero-field-splitting parameter (Δ), the temperature at which magnetic ordering occurs (θ) along with the g-factor (g) for each system. Applying the model to each dataset by method of a non-linear least squares fitting resulted in the values displayed in Table 3. The resulting fits are plotted on top of the experimental data points in Figure 5, showing good agreement.



Figure 5. Susceptibility multiplied by the temperature as a function of temperature. Fit including interactions and zero-field splitting.

The fit parameters from equations [1] and [2] are presented in Table 3Table for all three samples, along with parameters for two compounds from literature. The numerical value of the zero-field splitting parameter obtained for all three compounds are very similar, implying they experience very similar zero field splitting, as expected since the Dy^{III} ions in each respective compound are in very similar coordination environments.

 Table 3. Results of the non-linear least squares fitting applied to the model developed by McPherson et al.

		A		
Phase / Ref	Fitting	g -	Δa	θ ^ь
	Range (K)	factor		
Crystalline solid	1.8 - 400	1.25	1.3	-0.69
Liquid	1.8 - 400	1.22	1.45	-0.42
Glass	1.8 - 400	1.14	1.26	-0.88
Binuclear Mo ^v –Dy ^{III [5]}	1.8 - 400	1.34	92.8	+ 0.9
Mn ^{III} - Dy ^{III}	1.8 - 400	1.37	1.4	- 0.02K°
complexes ^[4]				

^aΔ = zero-field splitting parameter (cm⁻¹) x 10⁻³; ^bθ = Magnetic Ordering Temperature (K); ^c converted the wavenumber to Temperature for comparison.

It is the ligands that differ and give rise to the observed phase changes between compounds. The zero-field splitting parameters found for the samples here are of the same order of magnitude as those reported in literature.^[4,5] The magnetic ordering temperatures obtained for all compounds are under 1 K; this is as expected due to the localized nature of the 4f magnetism exhibited by RE ions. The Weiss parameters for all three compounds are identical, within uncertainty. This implies that the distance between the ions are the same, since the magnetic interaction decreases with distance between moments. Simply fitting using equation (1), ignoring the Weiss parameter for antiferromagnetic interaction, does not accurately reproduce the data (Figure S14 in the supplementary material). Hence, the decrease in χT at low temperatures must be a combination of both depopulation (zero field splitting) effects and antiferromagnetism. This holds for all the datasets and is consistent with magnetic dimers being present in all the measured compounds, with magnetic ordering temperatures under 1 K.

Herein, we demonstrate the design from crystalline to glass and most importantly, liquid dimeric Dy^{III} containing ionic compounds. Since speciation of metals in the liquid state can be a challenge, we have developed an innovative approach by using the magnetic properties to prove metal-metal interactions within the dimeric structure in the liquid phase. The same antiferromagnetic signature was observed in the crystalline, glass and liquid compounds a strong indicator that the magnetic dimers formed in the crystalline solid were in fact formed in the glass and liquid state as well.

Experimental Section

Experimental details can be found in the supplementary information.

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Entry for the Table of Contents

COMMUNICATION



Variation in anion and cation structure upon changing from crystalline solid to glass to liquid dimeric lanthanide-containing compounds.

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Designing Dimeric Lanthanide(III)-Containing Ionic Liquids

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1 Experimental Procedures

1.1 Materials and instrumentation.

 $Dy(CH_3COO)_34H_2O$, chemical grade, 99.9% and sodium acetate anhydrous, 99% were purchased from Alfa Aesar. $[C_4Mim][CH_3COO]_* \ge 95\%$ was purchased from Sigma Aldrich. $[P_{666\ 14}][CI]$ was purchased from Cytec and used without further purification. Nuclear magnetic resonance (NMR) spectra were recorded on a Bruker 400 MHz ultrashield spectrometer. TGA and DSC were conducted on a TA Q5000 and Q2000 equipped with a refrigerated cooling system 90 (temperature range -90 – 150 °C) and auto-sampler, respectively. In addition, elemental analysis was determined using a PerkinElmer 2400 series 2 CHNS/O analyser. FTIR was carried out on a PerkinElmer Spectrum 100 spectrometer and UV/Vis spectra were recorded on an Agilent, Cary 60 spectrophotometer.

1.2 X-ray crystallographical analysis.

Crystal data collection was performed on a Rigaku Single Crystal X-ray diffractometer, equipped with cryostat. Measurements were carried out at 100 K using a Cu K α , monochromator, λ = 1.54184 Å. Using Olex2, the structure was solved with the ShelXT structure solution program using Intrinsic Phasing and further refined with the ShelXL refinement package using least squares minimisation.

1.3 D.C. superconducting quantum interference device magnetometry.

The DC magnetic susceptibility of samples presented in this paper were measured using a Quantum Design MPMS XL SQUID, utilizing the reciprocal sample option (RSO). The magnetization was measured over a temperature range of 1.8 to 400 K. All samples were zero-field cooled from 400 K down to 1.8 K, and upon reaching the lowest temperature a field of 500 Oe was applied. The sample was then heated from 1.8 to 400 K in the applied field. Measurements of the magnetization were taken in 0.5 K intervals (when T < 20 K) and 3 K intervals (when T > 20 K) and the temperature was allowed to stabilize prior to any measurements being made. For verification of the paramagnetic behaviour (non-hysteretic behaviour) of the samples, data were collected upon cooling and warming and no differences were observed between the two protocols. Diamagnetic corrections were calculated for the relevant chemical structures using Pascal's constants and applied to all data.^[1] The measurement of the magnetic moment in our MPMS SQUID has an uncertainty of 5 % which arises mainly from differences in sample mounting and cannot be eliminated.

1.4 Synthesis of [C₄Mim]₂[Dy₂(CH₃COO)₈].

To a GC sample vial $[C_4Mim][CH_3COO]$ (59 mg, 0.297 mmol) and $Dy(CH_3COO)_3 \cdot 4H_2O$ (122.5 mg, 0.297 mmol) were added. The vial was sealed and wrapped in aluminium foil and placed in a Reacti-Therm to be slowly heated to 120°C. The vial was subsequently cooled over a period of 48 h to allow $[C_4Mim]_2[Dy_2(CH_3COO)_8]$ crystals to form. Anal. Calcd for $C_{32}H_{54}O_{16}N_4Dy_2$: C, 35.73; H, 5.06; N, 5.21. Found: C, 35.37; H, 4.99; N, 5.81.

1.5 Synthesis of $[P_{666 \ 14}]_2[Dy_2(CH_3COO)_8]$.

 $[P_{666\ 14}][CH_3COO]$ was prepared as follows; To a solution of $[P_{666\ 14}][CI]$ (5.1 g, 9.82 mmol) in ethanol, Na(CH_3COO) (0.81 g, 9.82 mmol) dissolved in ethanol was added. The reaction was left to stir for 48 hours at room temperature. Büchner filtration was used to remove the by-product NaCl. The ethanol was removed by rotary evaporation at 60 °C. Centrifugation was necessary to further remove NaCl. The final product was dried under vacuum at 70 °C. A pale yellow viscous liquid was obtained. ¹H NMR (Supplementary Fig. S2) (400 MHz, CDCl_3): δ = 2.43 (m, 8 H), 1.94 (s, 3 H), 1.47 – 1.24 (m, 48 H), 0.90 – 0.85 (m, 12 H) ppm. ¹³C NMR (Supplementary Fig. S3) (100 MHz, CDCl_3): δ = 176.86 (COO), 31.91 (CH₂COO), 30.90 – 21.88 (CH₂), 19.17 (CH₃), 14.10 (PCH₂) ppm. ³¹P NMR (Supplementary Fig. S4) (162 MHz, CDCl_3): δ = 33.02 ppm.

In a sealed vial [$P_{666\ 14}$][CH₃COO] (166 mg, 0.31 mmol) was then added to Dy(CH₃COO)₃·4H₂O (129.22 mg, 0.31 mmol) and the reaction mixture was stirred for 24 hours at 80 °C. The product was obtained as a pale-yellow glass.

1.6 Synthesis of [P_{666 14}]₂[Dy₂(C₇H₁₅COO)₈].

 $[P_{666 14}][C_7H_{15}COO]$ and $Dy(C_7H_{15}COO)_3$; xH₂O were prepared according to previously reported literature^{[2],[3]} described as follows; octanoic acid (5.53 g, 38.4 mmol) and [P_{666 14}]CI (19.94 g, 38.4 mmol) were dissolved in 75 mL of hexane. A solution of NaOH (1.54 g, 38.4 mmol, in 75 mL of H₂O) was added dropwise at room temperature (RT). The mixture was stirred at RT overnight. The upper organic phase was separated and washed four times with H₂O. Solvent was removed by rotary evaporator and the product was dried at 70°C under vacuum for 24 hours to yield [$P_{666 14}$][$C_7H_{15}COO$] as a viscous liquid. Dy($C_7H_{15}COO$)₃ was synthesised according to previously reported syntheses by Binnemans et al.[3] Octanoic acid was dissolved in an ethanol/water mixture and transformed to sodium octanoate by adding an equivalent amount of 1 M/L NaOH standard aqueous solution (addition dropwise via a burette). Dy (III) nitrate hydrate was dissolved in 100 mL ethanol/water mixture and added dropwise to the sodium octanoate solution. Upon addition a white precipitate formed immediately. The reaction was left to stir for a further 1 hour after which the white precipitate was filtered and washed with ethanol/water mixture and the product was dried overnight at 30°C. Dy octanoate was obtained as a fine white powder. $[P_{666 14}]_2[Dy_2(C_7H_{15}COO)_8]$: In a sealed vial $[P_{666 14}]_2[C_7H_{15}COO]$ (166 mg, 0.26 mmol) was added to Dy(C₇H₁₅COO)₃ (156.74 mg, 0.26 mmol) Under inert conditions to minimise the influence of water and the reaction mixture was stirred for 24 hours in ethanol at 70 °C. [P_{666 14}]₂[Dy₂(C₇H₁₅COO)₈] was obtained as a transparent pale viscous liquid. [P_{666 14}][C₇H₁₅COO] ¹H NMR (Supplementary Fig. S5 (400 MHz, CDCl₃) δ 2.44 (m, J = 14.3 Hz, 8 H), 2.15 (t, J = 15.6 Hz, 2 H), 1.58 – 1.45 (m, 18 H), 1.32 – 1.21 (m, 40 H), 0.86 – 0.80 (m, 15 H) ppm. ¹³C NMR (Supplementary Fig. S6) (100 MHz, CDCl₃): δ = 180.40 (COO), 39.30 (CH₂COO), 31.94 - 21.91 (CH₂), 19.15 (CH₃), 14.07 (P-CH₂) ppm. ³¹P NMR (Supplementary Fig. S7) (162 MHz, CDCl₃): δ = 32.98 ppm. Anal. Calcd for C₈₀H₁₆₀O₁₆P₂Dy₂: C, 63.05; H, 10.58. Found: C, 65.39; H, 11.46.

2 Results and Discussion

2.1 Powder X-ray Diffraction



Figure S1. Powder XRD pattern of the bulk [C₄Mim]₂[Dy₂(CH₃COO)₈] sample (black) and pattern simulated from single crystal structure (red).

2.2 NMR Spectra

Trihexyltetra(decyl)phosphonium acetate ([P_{666 14}][CH₃COO])











Figure S4. ³¹P NMR spectrum of [P_{666 14}][CH₃COO].





Figure S5. ¹H NMR spectrum of [P_{666 14}][C₇H₁₅COO].



Figure S6. ¹³C NMR spectrum of [P_{666 14}][C₇H₁₅COO].



2.3 Thermal gravimetric analysis.



Figure S8. TGA of $[C_4Mim]_2[Dy_2(CH_3COO)_8]$ (black dash), $[P_{666\ 14}]_2[Dy_2(CH_3COO)_8]$ (blue) and $[P_{666\ 14}]_2[Dy_2(C_7H_{15}COO)_8]$ (red).

2.4 Differential Scanning Calorimetry



Figure S9. DSC of [C₄Mim]₂[Dy₂(CH₃COO)₈]



Figure S10. DSC of [P_{666 14}]₂[Dy₂(CH₃COO)₈].



Figure S11. DSC of [P_{666 14}]2[Dy₂(C₇H₁₅COO)8].

Ln	Solid (°C)	Glass (°C)	Liquids (°C)
La	267 (T _{dec})	-60 (T _g), 250 (T _{dec})	
Ce			328 (T _{dec})
Nd	73 (T _m) 262 (T _{dec})	-78 (Tg),267 (T _{dec})	333 (T _{dec})
Sm	259 (T _{dec})	-75 (Tg), 244 (T _{dec})	341 (T _{dec})
Gd	260 (T _{dec})		317 (T _{dec})
Dy	118 (T _m), 292 (T _{dec})	-41.3 (Tg), 335 (T _{dec})	-80.40 (Tg), 348 (T _{dec})
Er	261 (T _{dec})		317 (T _{dec})
Y	261 (T _{dec})		322 (T _{dec})
Lu			-59.38 (Tg)

Table S1. Melting points, crystallisation temperatures, glass transitions and decompositiontemperatures of the crystalline, glassy and liquid series.

 T_m = melting point, T_g = glass transition temperature, T_{dec} = decomposition temperature.

2.5 UV/Vis Spectroscopy



Figure S12. UV/Vis spectrum of [P_{666 14}]₂[Dy₂(C₇H₁₅COO)₈]

Table S2. Assignment of the f - f transitions in $[P_{666 \ 14}]_2[Dy_2(C_7H_{15}COO)_8]$.

Wavelength (nm)	Transition
325.73	$^{6}\text{H}_{15/2} \rightarrow ^{6}\text{P}_{3/2}$
338.53	${}^{6}\text{H}_{15/2} \rightarrow ({}^{4}\text{F}{}^{4}\text{D})_{5/2}$
351.34	$^{6}H_{15/2} \rightarrow \ ^{6}P_{7/2}$
365.75	${}^{6}\text{H}_{15/2} \rightarrow {}^{6}\text{P}_{5/2}$
380.15	$^{6}H_{15/2} \rightarrow {}^{4}K_{17/2}$
388.15	${}^{6}\text{H}_{15/2} \rightarrow {}^{4}\text{I}_{13/2}$
399.36	$^{6}H_{15/2} \rightarrow {}^{4}F_{7/2}$
426.56	${}^{6}\text{H}_{15/2} \rightarrow {}^{4}\text{G}_{11/2}$
452.17	${}^{6}\text{H}_{15/2} \rightarrow {}^{4}\text{I}_{15/2}$
474.58	${}^{6}\text{H}_{15/2} \rightarrow {}^{4}\text{F}_{9/2}$
756.49	${}^{6}\text{H}_{15/2} \rightarrow {}^{6}\text{F}_{3/2}$
807.49	$^{6}\mathrm{H}_{15/2} \rightarrow ^{6}\mathrm{F}_{5/2}$

2.6 Crystallographic Data

Table S3. Crystal structure data.

Identification code	[C4Mim]2[Dy2(CH3COO)8]
Empirical formula	C ₃₂ H ₅₄ N ₄ O ₁₆ Dy ₂
Formula weight	1075.79 g/mol
Temperature/K	293(2)
Crystal system	triclinic
Space group	P-1
a/Å	8.4373(3)
b/Å	15.5572(4)
c/Å	15.9881(5)
α/°	97.705(2)
β/°	103.952(3)
γ/°	90.338(2)
Volume/Å ³	2016.71(11)
Z	29
ρ _{calc} g/cm ³	4.908
µ/mm ⁻¹	141.307
F(000)	2552.0
Crystal size/mm ³	
Radiation	CuKα (λ = 1.54184)
2O range for data collection/°	10.814 to 177.41
Index ranges	-10 ≤ h ≤ 10, -19 ≤ k ≤ 19, -18 ≤ l ≤ 20
Reflections collected	31451
Independent reflections	7713 [R _{int} = 0.0688, R _{sigma} = 0.0581]
Data/restraints/parameters	7713/0/499
Goodness-of-fit on F ²	1.089
Final R indexes [I>=2σ (I)]	R ₁ = 0.0476, wR ₂ = 0.1365
Final R indexes [all data]	R ₁ = 0.0557, wR ₂ = 0.1503
Largest diff. peak/hole / e Å ⁻³	1.21/-2.25

2.7 SQUID Magnetometry



Figure S13. Susceptibility and Inverse susceptibility as a function of temperature. Inset shows low temperature region of the inverse susceptibility in more detail.

The results obtained from fitting the inverse susceptibilities, as shown in Figure S13, using the Curie-Weiss law are summarized in Table S4. The Curie-Weiss law does not take any interactions or zero field splitting into account and therefore must be viewed cautiously. To avoid perturbing the obtained results, the fitting range was limited to 50 - 300 K, fitting to the low T region using the Curie-Weiss law is nonsensical as the compounds are not behaving as simple paramagnets below this temperature. Even at temperatures higher than 50 K one would expect the Stark levels to begin depopulating. However, all the extracted ordering temperatures (θ) are small and negative, implying deviation from paramagnetism at low T, this can be attributed to a range of effects as described in the main text, one of which may be antiferromagnetism.

Equation 1 presents the Curie-Weiss law that was used to fit to the inverse susceptibilities in the temperature range 50 to 300 K for all the compounds, the fitting was carried out using a linear least squares method and the fits were also weighted to account for the variation in error as a function of temperature.

$$\chi = \frac{C}{T-\theta}$$
, Equation 1

SUPPORTING INFORMATION Table S4 – Curie-Weiss fit results.

Phase	Fitting Range (K)	θ (K) – Ordering Temperature	Effective Moment Per Dy Ion (Bohr Magnetons)
Solid Crystal	50 - 300	-2.1	10.7
Liquid	50 - 300	-0.2	10.4
Glass	50 - 300	-1.6	9.7



Figure S14. Susceptibility multiplied by temperature, as a function of temperature. Fit only including zero field interactions. Image b) shows the low temperature region of image a) in greater detail.

The glassy sample returned an effective magnetic moment of 9.7 μ_B per Dy^{III} ion. This is 5% lower than the expected value; however, given that the measurement uncertainty is 5%, these values agree within uncertainty.

3 References

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4 Author Contributions

Éadaoin McCourt and Li Zhenyu carried out the experimental procedures and characterization of samples. Dr Kane Esien (KE) performed the SQUID magnetometery and data analysis of magnetic data. Dr Solveig Felton

(SF) and Dr Peter Nockemann conceived and supervised the study. All authors contributed to the reading and writing of the manuscript.