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In Situ Ellipsometric Monitoring of Gold Nanorod Metamaterials Growth

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ABSTRACT: An in situ transmission-based system has been designed to optically monitor the ellipsometry constants of a hyperbolic plasmonic metamaterial during electrochemical growth. The metamaterial, made from an array of vertically aligned gold nanorods, has demonstrated an unprecedented ability to manipulate the polarization of light using subwavelength thickness slabs, making in situ ellipsometric data a powerful tool in the controlled design of such components. In this work, we show practical proof-of-principle of this design method and rationalize the ellipsometric output on the basis of the modal properties of the nanorod metamaterial. The real-time optical monitoring setup provides excellent control and repeatability of nanostructure growth for the design of future ultrathin optical components. The performance of the ellipsometric method was also tested as a refractive index sensor. Monitoring refractive index changes near the metamaterial’s epsilon near zero (ENZ) frequency showed a sensitivity on the order of 500°/RIU in the ellipsometric phase for a metamaterial that shows 250 nm/RIU sensitivity in its extinction.

KEYWORDS: in situ, nanostructures, plasmonics, ellipsometry, hyperbolic, metamaterial, nanorods

1. INTRODUCTION

Optical in situ monitoring during growth and etching processes is a common procedure in the semiconductor industry. It is seen as a noninvasive method for controlling the fabrication and morphology of precision thin films. The most frequently used techniques include reflectance anisotropy spectroscopy (RAS), surface photon absorption (SPA), and ellipsometry. Although in situ optical monitoring of nanostructure fabrication is not unseen, their characterization is largely performed post growth and is often done at the expense of destroying the sample. Real-time in situ optical monitoring could open the door to the scalable and repeatable production of functional nanostructured systems, in particular resonant plasmonic nanostructures, whose optical properties are strongly sensitive to geometrical parameters, including refractive index, size, and shape. Real-time monitoring can also show the interplay and development of new modes as the growth of complex nanostructures progresses, providing a better understanding of their underlying behavior, but most of all allowing one to design optical functionalities on demand in vivo.

One of the most comprehensively studied plasmonic nanostructured systems, and also one of the most promising for practical implementation based on its ease and scalability to fabricate, are nanorod-based metamaterials. Nanorods of various materials can be grown in solution chemically or by electrodeposition into an alumina template. Gold nanorods are particularly interesting as they exhibit anisotropic optical properties throughout the visible and near-infrared spectral range, a property related to localized surface plasmon resonances (LSPR) across their long and short axes. When incorporated into an alumina template of suitable geometry, LSPR supported by nanorods when isolated, electromagnetically couple providing the formed metamaterial with unique optical properties that can be tailored to address a range of applications from biosensing to ultrafast optical switching.

It has recently been demonstrated that such a uniaxial nanorod metamaterial can be used as an ultrathin optical component to effectively manipulate the polarization of light. In particular, it was shown that both the reflection from and transmission through a $\lambda/20$ thick slab of the metamaterial can provide a 90° linear polarization rotation, a performance not observed in natural materials. The polarization conversion efficiency comes from the plasmonic nature of the modes supported by the metamaterial, providing a significant difference between the ordinary and extraordinary permittivities. The ability to manipulate the polarization of light is essential in many applications including liquid crystal displays, medical diagnostics, and telecommunications. For example, such hyperbolic metamaterials could replace conventional birefringent materials, such as rutile, which require tens
of micrometers to achieve similar polarization conversions. However, for nanorod metamaterial polarization converters to be adopted for practical use, precise engineering of the anisotropy in the metamaterial is required.

In this paper, we have devised a simple optical setup whereby the ellipsometric constants and the transmittance of the metamaterial are measured simultaneously during the electrochemical growth process. This in situ measurement technique not only improves upon the control and repeatability with which these metamaterials may be fabricated, but also provides a means for understanding their optical response as their modal structure evolves during growth. Additionally, given the importance of sensing in many applied fields, as in functional label-free biosensing, the sensitivity performance of the ellipsometric data to refractive index changes is further demonstrated.

2. METHODS

2.1. Nanorod Metamaterial Fabrication. The fabrication technique used to create gold (Au) nanorod metamaterials has been published elsewhere in great detail.15 The general geometry of the system is presented in Figure 1. Briefly, thin films of tantalum pentoxide (Ta₂O₅), gold, and aluminum are deposited using magnetron sputtering to form the sample template. The aluminum top layer is then anodized in a 0.3 M sulfuric acid bath at 35 V to produce an array of pores with an average center-to-center separation of \( d = 80 \text{ nm} \). In this instance the thickness of the aluminum layer was set to 400 nm. The aluminum is then chemically etched in a 30 mM NaOH solution to remove the barrier layer at the base of the pores. This step is also used to controllably widen the pores, thus determining the diameter, \( 2r \), of the nanorods electrochemically grown in the subsequent step. For the present study, the electrodeposition of gold was carried out at a constant voltage (−0.45 V) in a gold chloride electrolyte solution to grow nanorods of various lengths \( L \) in the pores. This final step completes the fabrication of the metamaterial (Figure 1b).

2.2. Metamaterial Structural Characterization. An experimental setup was designed to simultaneously measure the transmission and ellipsometric phase angles of the gold nanorods based metamaterial in situ during growth (Figure 2). The sample was mounted in the electrodeposition cell to probe both the sample’s extinction and ellipsometric parameters in transmission at an incidence angle of 45°. This configuration is chosen for diagnostics purposes only, to allow probing transmission and ellipsometric parameters simultaneously for the same angle of incidence on different paths. More generally, both

![Figure 1](image1.png)

![Figure 2](image2.png)
measured quantities can be obtained from ellipsometric measurements alone, therefore lifting the constraint on angle of incidence.

The blue outline in Figure 2 highlights the transmission measurement arm of the setup. The incident light, produced by a Tungsten halogen source in the visible and near-infrared regime (400–900 nm), is incident to the sample through a series of optical components controlling collimation and polarization. The latter was set to p-polarized light, where the electric field oscillates with a component both along the length and the diameter of the nanorods, i.e. in the (x, z) plane of Figure 1a and Figure 2, thus simultaneously probing both ordinary (x) and extraordinary (z) axis of the metamaterial. Transmission measurements recorded the spectrum of the light source after transit through the sample via an optical fiber bundle coupled to a spectrometer. The transmission and corresponding extinction were plotted during electrodeposition, giving an indication of the height of the nanorods and in turn providing direct control over their growth. The electrodeposition voltage and corresponding deposition current were also recorded across the working electrode and a platinum wire reference, giving additional information on the nanostructure growth characteristics.

2.3. In Situ Spectroscopic Ellipsometry. The red outlined arm in Figure 2 shows a schematic of the in situ ellipsometry measurement setup. For this proof-of-principle experiment, it is placed orthogonal to the transmission arm in order to provide both arms with the same angle of incidence. The ellipsometric measurements provide ellipsometric angles which are a measure of both amplitude ratio and phase shift between the p- and s-components of the light transmitted by the nanorod metamaterial, where the p- and s-components of the incident electric field are those in the (x, z) plane and along the y-axis, respectively (see Figures 1 and 2). Here we chose characterizing the ellipsometric angles in transmission as they provide for a direct insight into the hyperbolic properties of the metamaterial as will be shown below. When taken in transmission, the ellipsometric angles $\Psi$ and $\Delta$ are defined by the following equation.

Figure 3. In situ optical and ellipsometric response of the nanorod metamaterial measured during metamaterial growth. Average nanorod diameter and inter-rod spacing as determined by the template are 38 and 70 nm, respectively. The metamaterial growth is characterized by its thickness or equivalently the nanorod length. (a) Surface plot of the monitored extinction (−log10(Transmission)) spectra as a function of metamaterial thickness, assuming a linear relationship between growth time and metamaterial thickness, and (b) associated 2D cross-sectional plot. (g, h) Corresponding effective medium calculations. (c) Surface plot showing the spectral behavior of the ellipsometric amplitude angle ($\Psi_t$) as a function of metamaterial thickness during growth, and (d) associated 2D cross-sectional plot. (i, j) Corresponding effective medium calculations. (e) Surface plot showing the spectral behavior of the ellipsometric phase angle ($\Delta_t$) as a function of metamaterial thickness during growth, and (f) associated 2D cross-sectional plot. (k, l) Corresponding effective medium calculations. The general trend of the plotted data with increasing metamaterial thickness is indicated by arrows in the respective cross sections.
where \( t_p \) and \( t_s \) are the transmission coefficients for p- and s-polarized light, respectively, \( \tan \Psi = |E_p|/|E_s| \), with \( |E_p| \) and \( |E_s| \) the norm of the transmitted electric field for p- and s-polarized light, and \( \Delta \) is the relative phase difference between p- and s-polarized transmitted light. A commercial M-88 rotating analyzer ellipsometer (J.A. Woollam, Inc.) was used to perform the standard ellipsometry measurements. It was equipped with a Xenon arc lamp source for simultaneous measurements in the [450 nm-750 nm] wavelength range; a full spectrum was completed in 1/20 s and the WVASE software was used to compute the corresponding ellipsometric components \( \Psi \) and \( \Delta \) during nanorod growth. The instrument was mounted on an automated goniometer stage for measurements at variable angles of incidence, for postdeposition characterization.

2.4. Effective Medium Theory Modeling. The unique optical properties of gold nanorod metamaterials are ultimately governed by LSPR. Their optical response can be modeled using effective medium theories (EMTs), simplifying composite layers of nanoscale metallic inclusions with position and frequency-dependent permittivity \( \varepsilon_{eff}(\omega) \) embedded in a matrix material with position and frequency-dependent permittivity \( \varepsilon_{m}(\omega) \) into a layer with a frequency-dependent effective permittivity \( \varepsilon_{eff}(\omega) \) (Figure 1c, d). This effective permittivity can be found from the metamaterial’s volume fractions and polarizabilities within a Maxwell-Garnett approximation as shown before by R. Atkinson et al. In this formulation, the metamaterial layer is then considered a uniaxial anisotropic film (Figure 1d) with an effective permittivity expressed in the Cartesian reference system of Figure 1 as \( \varepsilon_{eff}(\omega) = \text{diag}(\varepsilon_x(\omega), \varepsilon_y(\omega), \varepsilon_z(\omega)) \), where \( z \) is taken along the nanorod length, and \( z \) is the extraordinary axis of the metamaterial, while the ordinary axis lies in the \((x, y)\) plane. Gold nanorod-based metamaterials are attracting a lot of attention due to their anisotropic optical response in the vicinity and beyond the so-called epsilon-near-zero (ENZ) frequency which marks the transition between elliptic and hyperbolic dispersion regimes. Hyperbolic materials are characterized by a negative effective permittivity \( \varepsilon_{eff}(\omega) < 0 \), where \( \varepsilon_{eff}(\omega) \) is the effective polarizability of the metamaterial as measured as a function of wavelength and metamaterial thickness, along with the corresponding ellipsometric parameters \( \Psi \) (Figure 3c, d) and \( \Delta \) (Figure 3e, f) obtained simultaneously during growth. The corresponding EMT calculation results are shown in Figure 3g–l, providing excellent agreement with the experimental observations, although it must be noted that the assumed metamaterial geometry during growth, including the average growth rate, can be at the origin of small discrepancies between experimental and EMT observations in Figure 3. In particular, the growth dynamics is a complex phenomenon and its assumed linear time behavior, especially at early stages, is an approximation that affects the comparison between experiments and calculations. Here we are making use of the EMT results to provide a general understanding of the experimental data and rationalize the ellipsometric response to allow their use in the design of ultrathin phase-shaping metamaterials with chosen and controlled polarization properties. As a result, approximations on electrochemical growth dynamics and metamaterial geometry during the growth do not affect the generality of our analysis below.

The overall extinction of the metamaterial increases steadily during growth as Au is deposited into the alumina pores to form the rods (Figure 3a, b). After about 100 s of growth time, or for a metamaterial thickness of about 200 nm, two distinct resonances can be observed, one at a wavelength of about 550 nm and another at around 625 nm. The short-wavelength peak, referred to as the T-mode in the literature, is associated with a transverse LSPR originating from free electrons oscillating across the diameter of the nanorods. The second peak, observed in the extinction at around 625 nm, corresponds to the resonant excitation of the free electron density along the long axis of the rods and is referred to as the L-mode in the literature. Further metamaterial growth leads to the observation of an increased spectral splitting of the two modes with a blue-shift of the T-mode to a wavelength of about 525 nm and a red-shift of the L-mode to about 650 nm. The spectral behavior of these resonances and their underlying nature as a function of metamaterial thickness has been thoroughly discussed elsewhere. The analysis of \( \Psi \) in Figure 3c, d is consistent with the measured extinction of Figure 3a, b. Indeed, for \( |E_p| > |E_s| \), we expect \( \Psi < 1 \), i.e., \( 45^\circ < \Psi < 90^\circ \), whereas for \( |E_p| < |E_s| \) we expect \( \Psi < 1 \), i.e., \( 0 \leq \Psi < 45^\circ \) with the case \( |E_p| = |E_s| \) corresponding to \( \Psi = 1 \) and \( \Psi = 45^\circ \). Figure 3c, d shows that the T-mode maximum corresponds to \( \Psi \approx 45^\circ \), which is expected for an angle of incidence of \( 45^\circ \), whereas \( \Psi \approx 0^\circ \) at the L-mode maximum with most of the transmitted field transmitted by the metamaterial in this spectral range being s-polarized. Again, this general behavior is also retrieved from the EMT considerations of Figure 3a–j. In fact, plotting Figure 4a such that \( \tan \Psi = |E_p|/|E_s| \approx \sqrt{T_p/T_s} \), where \( T_p \) and \( T_s \) are the metamaterial’s transmittance for p- and s-polarization, respectively, reproduces both the measured (Figure 3c) and EMT (Figures 3i and 4c) results both qualitatively and quantitatively very well, providing a direct link between \( \Psi \) and the resonant response of the metamaterial.
Although the $\Psi$ signal shows an absorptive shape, the phase signal $\Delta$ of Figure 3e, f shows a differential shape for metamaterial thicknesses exceeding ~150 nm. This thickness corresponds to the transition from the assembly from nanorods acting as an ensemble of weakly interacting resonators to a metamaterial made of strongly interacting nanorods. In fact, the differential shape is a signature of the dispersive behavior of the metamaterial transitioning between the elliptical and hyperbolic dispersion regimes. In the elliptic regime, the s-polarized field lags the p-polarized field (blue region in Figure 3e), while the opposite situation arises in the hyperbolic regime (red region in Figure 3e). Again, this observation is confirmed through the EMT calculations of Figure 3e, f and Figure 4c, d. Further insight can be provided simply by expressing the phase difference $\Delta \phi = 2 \pi \Delta (n_p L_p - n_s L_s)$, where $\lambda_0$ is the free-space wavelength of the transmitted light, $n_p$ and $n_s$ are the respective refractive indices experienced by p- and s-polarized waves, and $L_p$ and $L_s$ are the geometrical pathlengths for p- and s-polarized fields, respectively. Expressions for the refractive index are chiefly those governing the propagation of light in uniaxial materials such as $n_p = \sqrt{\varepsilon_p} = \sqrt{\varepsilon_0 + \sin^2 \theta (1 - \frac{\lambda_0}{\lambda})}$, where $\theta$ is the angle of incidence in air as measured from the z-axis. To a first approximation the pathlengths can be expressed as $L_{p,s} = \frac{d}{\cos \theta_{p,s}}$, where the angle of refraction for p- and s-polarized fields $\theta_{p,s}$ is found using Snell’s law of refraction $\sin \theta = \Re(\sqrt{\varepsilon_{p,s}}) \sin \theta_{p,s}$. The resulting map of $\Delta \phi$, plotted in Figure 4b, reproduces the main features of the experimental observations of Figure 3e and is in excellent qualitative agreement with the EMT calculations of Figure 3k and Figure 4d. The change in sign of the dephasing between p and s waves follows that of the term $\lambda_0 (n_p L_p - n_s L_s)$. In fact, for incident angles exceeding normal incidence Snell’s law predicts opposite trends for $n_p - n_s$ and $L_p - L_s$, with the geometric path length decreasing for increasing refractive index. As a result, $n_p - n_s$ and $L_p - L_s$ have a compensating effect on $\Delta$ but for the geometry considered here, the EMT calculations show that the behavior of $\Delta$ follows that of $n_p - n_s$ with max $(|n_p - n_s| > 1$, whereas max $(|L_p - L_s| \approx 50$ nm (see Figures S2–S4). For smaller metamaterial thicknesses, the length of the nanorods is close to their diameter leading to $n_p \approx n_s$ and subsequently $L_p \approx n_s$ and $\Delta \approx 0$ as measured in Figure 3c, d. As the length of the nanorods increases, the spectral behavior of $\Delta$ reflects that of $n_p - n_s$ with the s-polarized field leading the p-polarized field by more than $90^\circ$ in the spectral range of elliptic dispersion to the s-polarized field lagging the p-polarized field by almost $60^\circ$ in the spectral range of hyperbolic dispersion. The cross over where $\Delta = 0$ occurs at the L-mode resonance, slightly red-shifted from the ENZ frequency.

We finish our study by assessing the performance of the ellipsometric response of the nanorod metamaterial as a sensing platform. This is done in situ by monitoring the ellipsometric parameters in transmission while etching the alumina matrix surrounding the gold nanorod in the metamaterial. The change in the local refractive index resulting from this etch is a simple way of testing the refractive index sensitivity of ellipsometric measurements. These will be evaluated against the more commonly used resonant changes resulting from the same refractive index change measured simultaneously via the extinction of the metamaterial. For the experiment, the metamaterial is placed in a cell and immersed in 0.03 M sodium hydroxide solution with both ellipsometric parameters and extinction measured simultaneously as the alumina is gradually replaced by the etching solution changing the refractive index of the medium embedding the nanorods from $n \approx 1.6$ (alumina) to $n \approx 1.33$ (etching solution). With alumina as the embedding medium ($n \approx 1.6$), the metamaterial transitions from the elliptic to the hyperbolic regime at a wavelength of around 650 nm. This spectral range is the most sensitive to changes in refractive index when monitored in transmission as it corresponds to a resonance for incident TM-polarized light. The ellipsometric spectra recorded for changing refractive index for both $\Psi$ and $\Delta$ are shown in Figure 5. The spectral map for $\Psi$ (Figure 5a) and $\Delta$ (Figure 5b) show the total shift in the transmission wavelength of 70 nm from about 650 nm to about 580 nm. This corresponds to the full removal of the alumina template and replacement by the etching solution as the embedding refractive index is lowered from $n \approx 1.6$ to $n \approx 1.33$. This is confirmed by the extinction measurements (not shown) and gives a bulk sensitivity of approximately 250 nm/RIU when monitoring the L-mode resonance. This refractive index sensitivity is typical for nanorod-based metamaterials when monitored in transmission, although improvements by an order of magnitude can be achieved via both thermal annealing and by shifting the transition wavelength to lower frequencies, the latter increasing both sensitivity and dynamic range.

Both $\Psi$ and $\Delta$ data in Figure 5a, b demonstrate a behavior similar to our observations in Figure 3. The etching process starts at a time of around 30s (Figure 5a, b), following a nonlinear behavior until completion. This nonlinear behavior reflects both the density profile of the alumina matrix and is indicative of the etching mechanism. In fact, the alumina matrix density increases gradually away from the pores containing the nanorods, corresponding also to a gradual increase in its refractive index. As a result, when subjected to the etching solution, the alumina is uniformly removed in the matrix along the length of the nanorods first, where the matrix
has the lowest density, creating a shell around the nanorods whose thickness grows outward toward neighboring nanorods as the etching process proceeds.\(^{30,31}\) A saturation in the resonance shift is reached after a time of around 120 s as a result of the spectral dispersion of the metamaterial in this spectral range, which makes the 250 nm/RIU sensitivity obtained from the extinction measurements a lower estimate value. Figure 5c, d shows selected cross-sections of ψ and Δ, respectively, for various wavelengths in the 70 nm bandwidth of the resonance shift. Although the change in the amplitude parameter Ψ is limited to about 20° by the change in \(\sqrt{r_p/r_s}\), as explained earlier the dephasing angle Δ varies by more than 150° during etching over the same spectral range as a result of the transition between elliptic and hyperbolic dispersions. This translates into a sensing sensitivity for Δ on the order of 500°/RIU, far exceeding the 250 nm/RIU sensitivity simultaneously obtained from the extinction of the same metamaterial. Importantly, it must be noted that unlike extinction, the amplitude in the variations of Δ, and Ψ to a lesser extent, are robust and relatively independent of bandwidth as long as some shift in the resonance occurs as a result of the refractive index change to trigger the transition between the elliptic and hyperbolic dispersion regimes. By replacing the sample cell with a flow-cell one could create a fully functional label-free biosensor, which dispersions regimes. By replacing the sample cell with a flow-cell one could create a fully functional label-free biosensor, which.

**CONCLUSION**

We studied the in situ ellipsometric response of a nanorod-based plasmonic metamaterial during growth. The ellipsometric angle data, measured in transmission, were rationalized with respect of the optical properties of the metamaterial on the basis of the extinction measurements, measured simultaneously. The ellipsometric angles show significant variations in the vicinity of the ENZ frequency, where the metamaterial transitions from the elliptic to the hyperbolic dispersion regimes. In particular, the amplitude ratio shows an absorptive behavior with a maximum absolute change on the order of 50°, whereas the relative phase shows a dispersive behavior about the ENZ frequency with an amplitude change exceeding 150°. These variations have been associated with both the relative amplitude and relative phase of p- and s-waves transmitted through the metamaterial. The in situ ellipsometric data provide for a reliable method to grow these subwavelength-thick metamaterials to provide adjustable optical properties for applications as ultrathin polarization converters. Additionally, we assessed the sensitivity of the ellipsometric parameters to changes in the refractive index within the metamaterial. Here changes in the relative phase exceed 500°/RIU for a metamaterial that shows a 250 nm/RIU sensitivity in its extinction.

**ASSOCIATED CONTENT**

Supporting Information

The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/acsami.7b04129.

Figures S1–S4 (PDF)

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Notes

The authors declare no competing financial interest.

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