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A Stochastic Geometric Analysis of Device-to-Device Communications Operating Over Generalized Fading Channels

Young Jin Chun, Member, IEEE, Simon L. Cotton, Senior Member, IEEE, Harpreet S. Dhillon, Member, IEEE, Ali Ghrayeb, Senior Member, IEEE, and Mazen O. Hasna, Senior Member, IEEE

Abstract—Device-to-device (D2D) communications are now considered an integral part of future 5G networks, which will enable direct communication between user equipments and achieve higher throughputs than conventional cellular networks, but with the increased potential for co-channel interference. The physical channels, which constitute D2D communications, can be expected to be complex in nature, experiencing both line-of-sight (LOS) and non-LOS conditions across closely located D2D pairs. In addition to this, given the diverse range of operating environments, they may also be subject to clustering of the scattered multipath contribution, i.e., propagation characteristics which are quite dissimilar to conventional Rayleigh fading environments. To address these challenges, we consider two recently proposed generalized fading models, namely $\kappa$-$\mu$ and $\eta$-$\mu$, to characterize the fading behavior in D2D communications. Together, these models encompass many of the most widely utilized fading models in the literature such as Rayleigh, Rice (Nakagami-$m$), Nakagami-$m$, Hoyt (Nakagami-$q$), and One-sided Gaussian. Using stochastic geometry, we evaluate the spectral efficiency and outage probability of D2D networks under generalized fading conditions and present new insights into the tradeoffs between the reliability, rate, and mode selection. Through numerical evaluations, we also investigate the performance gains of D2D networks and demonstrate their superiority over traditional cellular networks.

Index Terms—5G, device-to-device network, $\eta$-$\mu$ fading, $\kappa$-$\mu$ fading, rate-reliability trade-off, stochastic geometry.

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I. INTRODUCTION

A. Related Works

The recent unprecedented growth in mobile traffic has compelled the telecommunications industry to come up with new and innovative ways to improve cellular network performance to meet the ever increasing data demands. This has led to the introduction of fifth generation (5G) networks which are expected to provide 1000 fold gains in capacity while achieving latencies of less than 1 millisecond [1]. Device-to-device communications are a strong contender for 5G networks [2] that allow direct communication between user equipments (UEs) without unnecessary routing of traffic through the network infrastructure, resulting in shorter transmission distances and improved data rates than traditional cellular networks [3].

Currently, D2D communication is standardized by the 3rd Generation Partnership Project (3GPP) in LTE Release 12 to provide proximity based services and public safety applications [4]. In parallel to the standardization efforts, D2D communications have been actively studied by the research community. For example, in [5], the authors have proposed D2D as a multi-hop scheme, while in [6] and [7], the work conducted in [5] has been extended to demonstrate that D2D communications can improve spectral efficiency and the coverage of conventional cellular networks. Additionally, D2D has also been applied to multi-cast scenarios [8], machine-to-machine (M2M) communications [9], and cellular off-loading [10].

While D2D communications offer many advantages, they also come with numerous challenges. These include the difficulties in accurately modeling the interference induced by cellular and D2D UEs, and consequently optimizing the resource allocation based on the interference model. Most of the previous works published in this area have relied on system-level simulations with a large parameter set [11], meaning that it is difficult to draw general conclusions. Recently, stochastic geometry has received considerable attention as a useful mathematical tool for interference modeling. Specifically, stochastic geometry treats the locations of the interferers as points distributed according to a spatial point process [12]. Such an approach captures the topological randomness in the network geometry, offers high analytical flexibility and achieves accurate performance evaluation [13]–[17].
Much work has also been done on evaluating the performance of D2D networks over Rayleigh fading channels. In [18], the authors have compared two D2D spectrum sharing schemes (overlay and underlay) and evaluated the average achievable rate for each scheme based on the stochastic geometric framework. In [19], the authors extended the work conducted in [18] by considering a D2D link whose length depends on the user density. In [20], the authors proposed a flexible mode selection scheme which makes use of truncated channel inversion based power control for underlaid D2D networks. Notwithstanding these advances, limited work has been conducted to consider D2D networks with general fading channels, for example in [21], the authors have considered underlaid D2D networks over Rician fading channels and evaluated the success probability and average achievable rate.

**B. Motivation and Contributions**

In 5G networks and especially for D2D communications, fading environments will range from homogeneous and circularly symmetric through to non-homogeneous and non-circularly symmetric. For example, the METIS project has already demonstrated that the physical channels of 5G networks can be inhomogeneous with clusters of non-circularly symmetric scattered waves [22]. Clearly in this case, the assumption of traditional, homogeneous, linear and single cluster fading models such as Rayleigh will no longer be sufficient and we must look towards other more general and realistic models such as $\kappa$-$\mu$ [23]–[25] and $\eta$-$\mu$ [23], [26].

Influenced by this, we consider the $\kappa$-$\mu$ fading model which accounts for homogeneous, linear environments with line-of-sight (LOS) components and multiple clusters of scattered signal contributions, while the $\eta$-$\mu$ fading model represents inhomogeneous, linear environments with non-line-of-sight (NLOS) conditions and multiple clusters of scattered signal contributions.

As discussed earlier, most of the existing work in stochastic geometry for wireless networks has been focused on Rayleigh fading environments, owing to its tractability and favorable analytical characteristics. The signal-to-noise-plus-interference ratio (SINR) distributions for general fading environments require evaluating the sum-products of aggregate interference where several approaches have been proposed to facilitate the derivation, most notably:

**1)** The conversion method based on displacement theorem was used in [27]–[30]. This method treats the channel randomness as a perturbation in the location of the transmitter and transforms the original network with arbitrary fading into an equivalent network without fading. Although the conversion method can be applied to any fading distribution, it is more tractable for handling large-scale shadowing effects. Specifically, if one applies the conversion method to small-scale fading, the resulting equivalent model will have no fading, thereby the Laplace transform-based approach can not be utilized.

**2)** The series representation method was used in [21] and [31]. This approach expresses the interference functionals as an infinite series of higher order derivative terms [32] given by the Laplace transform of the interference power. While the series representation method provides a tractable alternative for handling general fading, it often leads to situations where it is difficult to derive closed form expressions.

**3)** The integral transform based approach was used in [33]–[35], where either the Fourier transform (FT), Laplace transform (LT), characteristic function (CF) or moment generating function (MGF) is utilized. For instance, Gil-Pelaez’s inversion formula was used in [33] to find the distribution of the SINR using the MGF. However, Gil-Pelaez’s inversion formula involves an integral over the complex plane and the MGFs of the related random variables may not always exist. The Plancherel-Parseval theorem was used in [34] and [35] to calculate the expectation of an arbitrary function of the interference using the FT (or LT). Although the Plancherel-Parseval theorem provides a general framework, it often involves complex multi-fold integration and results in intractable expressions.

Motivated by these approaches and their limitations, we adopt a stochastic geometric framework to facilitate the performance evaluation of D2D networks over generalized fading channels; namely, $\kappa$-$\mu$ and $\eta$-$\mu$. We consider a D2D network overlaid upon a cellular network where the spatial locations of the mobile UEs as well as the base stations (BSs) are modeled as Poisson point processes (PPPs). The adopted framework can evaluate the average of an arbitrary function of the SINR, thereby enabling the estimation of the average rate and outage probability.

The main contributions of this paper may be summarized as follows.

1) We consider generalized fading conditions, namely, (i) $\kappa$-$\mu$ and (ii) $\eta$-$\mu$ fading, to account for various small-scale fading effects, such as LOS/NLOS conditions, multipath clustering, and power imbalance between the in-phase and quadrature signal components. These two models together encompass most of the popular fading models proposed in the literature. We utilize the series representation of the $\kappa$-$\mu$ and $\eta$-$\mu$ distributions to improve tractability and achieve closed form expressions.

2) We analyze the Laplace transform of the interference over $\kappa$-$\mu$ and $\eta$-$\mu$ fading channels and derive a closed form expression for the D2D and cellular links. By using a channel inversion based power control, we derive the Laplace transform of the interference in a closed form that does not involve an integral expression.

3) We exploit a novel stochastic geometric approach for evaluating the performance of D2D networks over generalized fading channels. This approach enables us to evaluate the average of an arbitrary function of the SINR as a closed form expression. We invoke the proposed stochastic geometric approach to evaluate the spectral efficiency and outage probability of D2D networks and compare that to the performance of conventional cellular networks. Furthermore, we study the trade-off among a number of performance metrics,
which can provide invaluable insights that may be used to optimize future network design.

The remainder of this paper is organized as follows. We describe the system model in Section II and the generalized fading models in Section III. We introduce the interference of cellular and D2D networks in Section IV, then utilize a stochastic geometric approach to evaluate the spectral efficiency and the outage probability of D2D networks in Section V. We present numerical results in Section VI and conclude the paper in Section VII with some closing remarks.

II. SYSTEM MODEL

A. Network Model

We consider a D2D network overlaid upon an uplink cellular network where a UE can directly communicate with other UEs without relying on the cellular infrastructure if a certain criterion is met. The overlaid spectrum access scheme allocates orthogonal time/frequency resources to the cellular and D2D transmitters by dividing the uplink spectrum into two non-overlapping portions. The overlay D2D network excludes cross-mode interference between cellular and D2D UEs and achieves a reliable link quality at the cost of lower spectrum utilization. Specifically, a fraction $\beta$ of the spectrum is assigned to D2D communications and the remaining $1 - \beta$ is allocated to cellular communications, where $0 < \beta < 1$.

Fig. 1 depicts a high level overview of the system model where the locations of the nodes are modeled as a spatial point process in $\mathbb{R}^2$. The UEs are assumed to form a homogeneous PPP $\Phi = \{X_i\}$ with intensity $\lambda$ and each UE $X_i$ has associated parameters that collectively form a marked PPP $\tilde{\Phi}$ as follows

$$\tilde{\Phi} = \{(X_i, q_i, L_i, P_i)\}, \quad (1)$$

where $L_i$ is the distance between the $i$-th UE and its intended receiver (referred to henceforth as the link distance), and $P_i$ is the transmit power of the $i$-th UE. The parameter $q_i$ indicates the inherent type of the $i$-th transmit UE which may be a potential D2D UE with probability $q = P(q_i = 1)$, or a cellular UE with probability $1 - q$, where $q \in [0, 1]$. For notational simplicity, we denote by $L_{d}$ the link length between a typical cellular UE and the associated BS. Similarly, $L_{d}$ represents the link length between a typical D2D UE and the D2D receiver UE. The receiver can be either a cellular BS or D2D receiver UE depending on the associated UE type. The cellular BSs are assumed to be uniformly distributed as PPP $\Psi$ with intensity $\lambda_b$. The D2D receiver UEs are randomly distributed around their associated D2D UE according to the distribution of the link length $L_{d}$, which is described later in (6). The performance analysis is performed for the typical receiver, which is assumed to be located at the origin due to the stationarity of this setup. The notations used in this paper are summarized in Table I.

B. Mode Selection and UE Classification

The operating mode of the UE $X_i \in \Phi$ is determined by two factors; 1) the inherent type ($q_i$) and 2) the mode selection policy. If $q_i = 0$, then the UE $X_i$ is a cellular UE and always connects to its closest BS (which is equivalent to the so-called maximum average received power association). If $q_i = 1$, then $X_i$ is a potential D2D UE which may use either cellular or D2D mode based on the adopted mode selection policy.

We assume a distance-based mode selection scheme [18]. That is, a potential D2D UE chooses the D2D mode if the D2D link length is smaller than or equal to a predefined mode selection threshold $\theta$, i.e., $L_{d} \leq \theta$. Otherwise, cellular mode is selected. Therefore, the complete set of transmit UEs $\Phi$ can be divided into two spatial point processes as follows

- UEs operating in cellular mode:
  $$\Phi_c \text{ with intensity } \lambda_c = [(1 - q) + qP(L_{d} > \theta)] \lambda, \quad (2)$$

- UEs operating in D2D mode:
  $$\Phi_d \text{ with intensity } \lambda_d = qP(L_{d} \leq \theta) \lambda. \quad (3)$$

1) Cellular Mode: We assume full-buffer transmission and orthogonal multiple access in the cellular uplink implying that at most one transmitter is active per cell over a given resource block. The locations of the active UEs (in the cellular mode) scheduled over the same resource block as the typical receiver are assumed to follow the point process $\Phi^c \subset \Phi_c$.

Due to the restriction that at most one point of $\Phi^c$ can lie in each cell of the Poisson Voronoi tessellation formed by $\Psi$, it is not straightforward to characterize $\Phi^c$. This is one of the key reasons why the exact uplink analysis for this setup has not yet been performed. Interested readers are advised to refer to [36] for more detailed discussion on user

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![Diagram](image)
distributions in cellular networks. For high user densities, one generative model that closely emulates this scenario while also being quite realistic from the actual system perspective is to assume that there is exactly one point of \( \Phi^e_c \) in each cell of the Poisson-Voronoi tessellation formed by \( \Psi \) (also assumed to be uniformly distributed in that cell). Viewed this way, \( \Phi^e_c \) can be thought of as a Poisson-Voronoi perturbed lattice [37]. However, due to the correlation between \( \Phi^e_c \) and \( \Psi \) (which correlates various link distances), the exact analysis for even this setup is also not known. That being said, it is possible to capture some of this correlation by approximating \( \Phi^e_c \) with a non-homogeneous PPP \( \hat{\Phi}_c \) with a distance-dependent intensity function. This uplink model was introduced and exploited in [13] and [38]–[40] and has been shown to provide a reasonable approximation for the SINR distribution of the cellular link. Please refer to [13] and [38] for more details about the uplink problem and this approximation approach.

Assumption 1: Given that the typical receiver is located at the origin, the set of active interfering UEs \( \Phi^e_c \) in the cellular mode that are scheduled within the same resource block is approximated by a non-homogeneous PPP \( \hat{\Phi}_c \) with intensity

\[
\hat{\lambda}_b = \frac{\lambda_b}{\pi d^2} 
\]

where \( \lambda_b = \frac{10}{\pi 500} \) and \( \xi \) is a fitting parameter that affects the average D2D link distance. Specifically, the scale parameter

\[
\lambda = \frac{10}{\pi 500} 
\]

2) D2D Mode: We model the D2D link length \( L_d \) using a Rayleigh distribution [18]

\[
f_{L_d}(r) = \frac{2\pi \hat{\lambda}_b r \exp\left(-\frac{\lambda_b \pi r^2}{\xi}\right)}{1 - \exp\left(-\frac{\lambda_b \pi d^2}{\xi}\right)}, \quad 0 \leq r \leq d, 
\]

where \( L_c \) is the link distance between the typical cellular UE and the connected BS at the origin, \( L_i \) represents the link distance between the interfering cellular UE \( X_i \) and its associated BS, and \( \| X_i \| \) is the distance between the interfering UE \( X_i \) and the BS located at the origin.\(^2\)

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\]

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\]

where \( \hat{\lambda}_b = \frac{\lambda_b}{\pi d^2} \) and \( \xi \) is a fitting parameter that affects the average D2D link distance. Specifically, the scale parameter

\[
\lambda = \frac{10}{\pi 500} 
\]
of the Rayleigh distribution in (6) is \( \sigma = \sqrt{\frac{4}{\pi e}} \) and the average D2D link distance is given by \( \mathbb{E}[L_d] = \sigma \sqrt{\frac{4}{\pi e}} \). D2D mode utilizes ALOHA with transmit probability \( \varepsilon \) on each time slot, where \( 0 \leq \varepsilon \leq 1 \). Since the potential D2D UEs in D2D mode follow a location independent thinning process [18], the set of UEs operating in the D2D mode is distributed according to a homogeneous PPP \( \Phi_d \) with intensity \( \lambda_d = \rho P(L_d \leq \theta)\lambda \), which is independent to the set of UEs in the cellular mode. Similarly, due to the location independent thinning induced by the ALOHA scheme, the set of active interfering UEs that gain access to the channel resource is distributed as a homogeneous PPP \( \varepsilon \Phi_d \) with intensity \( \varepsilon \lambda_d \), where \( \varepsilon \Phi_d \) is a subset of \( \Phi_d \), i.e., \( \varepsilon \Phi_d \subset \Phi_d \).

**Assumption 2:** In order to maintain tractability, we assume that \( \varepsilon \Phi_d \) and \( \Phi_c \) are independent.

### C. Channel Inversion-Based Power Control

The received power at the origin from the UE \( X_i \) is \( W = P_t \tau \|X_i\|^{-\alpha} G_i \), where \( P_t \) is the transmit power of the UE \( X_i \), \( \|X_i\| \) is the distance from \( X_i \) to the origin, \( \alpha \) is the path-loss exponent, \( \tau \) is the path-loss intercept at unit distance \( \|X_i\| = 1 \), and \( G_i \) represents the small-scale fading. The coefficients \( \{G_i\} \) of each link are assumed to be independent of one another.

We assume channel inversion based power control, i.e., \( P_t = L_t^{\mu} \). Then, the received power is \( W = \tau G_i (L_t/\|X_i\|)^\mu \) for an interference link and \( W = \tau G_i \) for the intended link. The transmit power of the UEs operating in the cellular mode \( X_i \in \Phi_c \) is \( P_d = L_t^{\kappa} \), whereas that of the potential D2D UEs in the D2D mode \( X_i \in \Phi_d \) is \( P_d = L_t^{\mu} \) given that \( L_d \leq \theta \). Since a potential D2D UE may use either a D2D mode or a cellular mode, for the purpose of our calculations, its transmit power can be interpreted as the weighted average of the two operating mode events, i.e., \( P_d = P(L_d \leq \theta)P_d + P(L_d > \theta)P_c \). Higher order moments of the transmit power for each mode are evaluated in the following lemma.

**Lemma 1:** The \( l \)-th moments of the transmit power of a cellular UE (\( P_c \)), a potential D2D UE in D2D mode (\( P_d \)), and a potential D2D UE (\( P_d \)) are respectively given by

\[
\begin{align*}
\mathbb{E}[P_c^l] &= \frac{\Gamma \left( \frac{l}{\delta} + 1 \right)}{(\lambda_b \pi)^{\frac{l}{\delta}}}, \\
\mathbb{E}[P_d^l] &= \frac{1}{(\lambda' \pi)^{\frac{l}{\delta}}} \left[ \frac{\gamma \left( \frac{l}{\delta} + 1, \frac{l'}{\delta} \theta^2 \right)}{1 - \exp \left( -\frac{l'}{\delta} \theta^2 \right)} \right], \\
\mathbb{E}[P_d^l] &= \frac{\exp \left( -\frac{l'}{\delta} \theta^2 \right)}{(\lambda' \pi)^{\frac{l}{\delta}}} \frac{\Gamma \left( \frac{l}{\delta} + 1 \right)}{(\lambda_b \pi)^{\frac{l}{\delta}}} + \frac{1}{(\lambda' \pi)^{\frac{l}{\delta}}} \frac{\gamma \left( \frac{l}{\delta} + 1, \frac{l'}{\delta} \theta^2 \right)}{1 - \exp \left( -\frac{l'}{\delta} \theta^2 \right)},
\end{align*}
\]

where \( \delta = \frac{\alpha}{2}, l > 0 \) is a positive real-valued constant, \( \lambda' = \frac{\alpha}{\pi} \), \( \theta \) is the mode selection threshold, \( \Gamma(t) \) is the gamma function.

3We have isolated and focused on studying the impact of the small scale fading upon the system model proposed here. Nonetheless, the model can be readily adapted to include shadowing by using the approach in [29, Lemma 1].
The PDF, l-th moment and Laplace transform of $G$ are respectively given by [23], [24], [41]

$$f_G(x) = \frac{\mu K_{\eta}^{\frac{-1}{\eta}}}{\kappa^{\frac{-1}{\eta}} \mu^{\frac{2}{\eta}} \Omega_{\eta \mu}^{\frac{1}{\eta}} \Gamma(\mu)} e^{-\frac{\kappa}{\mu} \Omega_{\eta \mu}} I_{\mu - 1} \left(2 \sqrt{\frac{\kappa}{\mu} \Omega_{\eta \mu}} x \right),$$

$$E\left[G^l\right] = \frac{\Omega_{\eta \mu}^{\frac{-1}{\eta}} \Gamma(\mu + 1)}{\mu^{\frac{2}{\eta}} \Gamma(\mu)} F_1(\mu + l; \mu; \mu \kappa),$$

$$L_G(s) = E\left[\exp(-sG)\right] = \left(1 + s \Omega_{\eta \mu}\right)^{-\mu} \exp(-\frac{\mu \kappa}{1 + s \Omega_{\eta \mu}}),$$

where $\tilde{w} = E[G]$, $\kappa$, $\mu$ and $l$ are positive real values, $\Omega_{\eta \mu} \triangleq \tilde{w}^{\frac{1}{\eta}} \mu^{\frac{2}{\eta}}$, $I_{\nu}(x)$ represents the modified Bessel function of the first kind, and $F_1(a; b; c; x)$ is the confluent hypergeometric function.

2) The $\eta-\mu$ Distribution: The $\eta-\mu$ distribution is used to represent small scale fading under non-line-of-sight (NLOS) conditions in inhomogeneous, linear, non-circularly symmetric environments [23], [26]. It is a general fading distribution that includes Hoyt (Nakagami-$q$), One-sided Gaussian, Rayleigh, and Nakagami-m as special cases (See Table II).

The received signal in an $\eta-\mu$ fading channel is composed of clusters of multipath waves. The in-phase and quadrature components of the fading signal within each cluster are assumed to be either independent with unequal powers or correlated with identical powers. The parameter $\eta$ denotes the scattered-wave power ratio between the in-phase and quadrature components and $2\mu$ represents the real valued extension of the number of multipath clusters.

The PDF, l-th moment and Laplace transform of $G$ are respectively given by [23], [24], [41]

$$f_G(x) = \frac{2 \sqrt{\pi} \mu^{rac{1}{2}} x^{rac{1}{2}} \mu}{\Gamma(\mu) H^{rac{1}{2}} \mu^{rac{1}{2}} \xi^{rac{1}{2}} \Omega_{\eta \mu}^{\frac{1}{2}}} e^{-\frac{\kappa}{\mu} \Omega_{\eta \mu}} I_{\mu - 1} \left(\frac{x}{\Omega_{\eta \mu}}\right),$$

$$E\left[G^l\right] = \frac{\Omega_{\eta \mu}^{\frac{-1}{\mu}} \Gamma(2\mu + l)}{\mu^{\frac{1}{2}} \Gamma(2\mu)} \times 2 F_1\left(\mu + \frac{l}{2}; \frac{1}{2}; \mu + \frac{1}{2}; \frac{(H^2)}{4}\right),$$

$$L_G(s) = \frac{1}{\mu H} \left[\left(1 + s \Omega_{\eta \mu}\right)^2 - \left(\frac{H^2}{4}\right)^2\right]^{-\mu},$$

where $\tilde{w} = E[G]$, $\eta$, $\mu$ and $l$ are positive real values,

5) $h = \frac{2 + \eta + 1}{\eta}, H = \frac{\eta - 1}{\eta}, \Omega_{\eta \mu} \triangleq \tilde{w}^{\frac{1}{\eta}} \mu^{\frac{2}{\eta}}$ and $2 F_1(a; b; c; x)$ is the Gaussian hypergeometric function.

IV. INTERFERENCE MODEL OF THE OVERLAY D2D NETWORK

In this section, we introduce the interference of cellular and D2D links and derive the Laplace transform of the interference for the generalized fading channels considered here.

A. D2D Mode

Let us consider a D2D link, where co-channel interference is generated by potential D2D UEs operating in D2D mode. Based on (8), the effective interference at the intended D2D receiver is

$$I_d = \sum_{X_j \in \Phi_d \setminus \{X_0\}} G_j L_j^d ||X_j||^{-\alpha}.$$

The point process $\Phi_d$ in (11) is similar to the D2D process in [18] since both models assume distance-based mode selection as well as Rayleigh distributed D2D link lengths. Nonetheless, the fading model in (11) affects the distribution of the aggregate interference and the corresponding distribution parameters, which is characterized by the Laplace transform of $I_d$ as follows.

**Lemma 2:** For overlay D2D, the Laplace transform of the interference at the D2D receiver is

$$L_{I_d}(s) = \exp(-c_d \cdot s^\delta),$$

where the constant terms $c_d$ for the $\kappa-\mu$ and $\eta-\mu$ fading distributions are respectively given by

$$\kappa - \mu : \frac{c_0 \Omega_{\eta \mu}^{\frac{-1}{\mu}} \Gamma(\mu + \delta - 1)}{\delta \mu^{\frac{1}{2}} \Gamma(\mu)} F_1\left(\mu, \delta; \mu; \mu \kappa\right),$$

$$\eta - \mu : \frac{c_0 \Omega_{\eta \mu}^{\frac{-1}{\mu}} \Gamma(2\mu + \delta - 1)}{\delta \mu^{\frac{1}{2}} \Gamma(2\mu)} \times 2 F_1\left(\mu + \frac{\delta + 1}{2}; \mu + \frac{\delta}{2}; \mu + \frac{1}{2}; \frac{H^2}{4}\right),$$

with $c_0 \triangleq q_{\mu \eta} \cdot \gamma(2, \lambda \pi \theta^2, \xi^{-1}), \delta = \frac{\zeta}{\tilde{w}}, \tilde{w} = E[G]$, the fitting parameter $\xi$ and ALOHA transmit probability $\epsilon$.

**Proof:** See Appendix III.

B. Cellular Mode

Based on Assumption 1, we model the interference at the cellular BS by a non-homogeneous PPP $\Phi_c$ with intensity $\lambda_c (1 - \exp(-\pi \lambda_b d^2))$ [13], [38], where the interference in (8) is

$$I_c = \sum_{X_j \in \Phi_c \setminus \{X_0\}} G_j L_j^c ||X_j||^{-\alpha},$$

and the Laplace transform of $I_c$ is evaluated as follows.

**Lemma 3:** For overlay D2D, the Laplace transform of the interference at the cellular BS is

$$L_{I_c}(s) = \exp(-s W(s)),$$

where $W(s)$ for each channel is respectively given by

$$W(s) = \frac{s \cdot \kappa - \mu \cdot \Omega_{\eta \mu}^{\frac{-1}{\mu}}}{(1 - \delta) e^\mu \frac{1}{\mu} \Gamma(\mu + 1, 1 - \delta; 2 - \delta, -s \Omega_{\eta \mu})} + (1 + s \Omega_{\eta \mu})^{-\mu} \exp(-\mu \kappa \frac{s \Omega_{\eta \mu}}{s \Omega_{\eta \mu} + 1},$$

for the $\kappa-\mu$ fading and

$$W(s) = \frac{2 \Omega_{\eta \mu}^{\frac{-1}{\mu}} H^2}{(1 - \delta) e^\mu \frac{1}{\mu} \Gamma(\mu + 1, 1 - \delta; 2 - \delta, -s \Omega_{\eta \mu})} \times 2 F_1\left(1, 1 - \delta - 2n; 2 - \delta; -s \Omega_{\eta \mu}\right) + s \Omega_{\eta \mu}^{\frac{-1}{\mu}} \exp\left(-\mu \kappa \frac{s \Omega_{\eta \mu}}{s \Omega_{\eta \mu} + 1},$$

for the $\eta-\mu$ fading and

$$+ h^{-\mu} \left(1 + s \Omega_{\eta \mu}\right)^2 - \left(\frac{H}{4}\right)^2 \right)^{-\mu} - 1.$$
for the $\eta$-$\mu$ fading, $\delta = \frac{2}{\alpha}, \Omega_{\kappa\mu} = \frac{\bar{\mu}}{\mu(1+\kappa)}$ and $\Omega_{\eta\mu} = \frac{\bar{\mu}}{2\mu^2}$.

**Proof:** See Appendix IV.

**Remark 1:** The Laplace transform of the interference in the cellular uplink is invariant to the node density $\lambda_b$. A similar behavior was observed in [38] for Rayleigh fading and Lemma 3 extends this invariance property to any fading model that can be represented by the $\kappa$-$\mu$ and $\eta$-$\mu$ distributions.

**Remark 2:** The invariance property and analytical tractability of Lemma 3 are special properties that only hold for the full channel-inversion based power control, i.e., $P_c = L^c_\kappa$. If we use fractional power control [13], [38], i.e., $P_c = L^c_\kappa$ with $0 \leq \epsilon \leq 1$, then the Laplace transform of the aggregate interference depends on the BS density and the integral in (48) can not be easily partitioned into two single integral expressions, which significantly complicates the derivation.

\[ E \left[ g \left( \frac{G_0}{I + N} \right) \right] = g(0) + \sum_{n=0}^{\infty} a_n \phi_{\eta\mu}(n), \]

where $\phi_{\eta\mu}(n)$ and $g_i(z)$ denote the following expressions

\[ \phi_{\eta\mu}(n) = \int_0^\infty g_{2\mu+2n}(z) L_\mu \left( \frac{z}{\Omega_{\eta\mu}} \right) \exp \left( -\frac{Nz}{\Omega_{\eta\mu}} \right) dz, \]

\[ g_i(z) = \frac{1}{\Gamma(2\mu + 2n)} dz \frac{\Gamma(z + 2\mu + 2n - 1)}{\Gamma(z + 2\mu + 2n)}, \]

**Proof:** See Appendix V.

Theorem 1 provides a general framework to evaluate arbitrary system performance measures when the received signal power $G_0$ follows either a Gamma distribution or a mixture of Gamma distributions, which includes $\kappa$-$\mu$ and $\eta$-$\mu$ and the majority of the most popular linear fading models utilized in the literature. We note that Theorem 1 makes no assumption on the underlying distribution of the constituent interference channels. Therefore it can be applied even when the intended signal and interfering links are described by different fading models.

In the following, we apply Theorem 1 and Lemmas 1-3 to evaluate various performance measures for overlaid D2D networks.

\[ R = \Delta \cdot E \left[ \log \left( 1 + \frac{G_0}{I + N} \right) \right]. \]

In D2D mode, due to the ALOHA medium access, $\epsilon$ percent of the transmitting UEs will gain access to the spectrum resource for the D2D transmission, which is $\beta$ fraction of the available spectrum. In cellular mode, $1 - \beta$ fraction of the available spectrum is allocated to the cellular transmission and only one uplink transmitter within each cell can stay active at any given resource block due to the orthogonal multiple access. Therefore, the amount of accessible spectrum resource is $\Delta_d = \beta \epsilon$ for the D2D transmission and $\Delta_c = (1 - \beta) E \left[ M \right]$ for the cellular transmission, where $M$ is the number of potential cellular UEs within a cell. The average $E \left[ M \right]$ is evaluated in [18] as $E \left[ M \right] = \frac{\lambda_c}{\lambda_e} \left( 1 - e^{-\frac{\lambda_c}{\lambda_e}} \right)$ where $\lambda_c = [(1 - \epsilon) + \epsilon P(L_d > \theta)] \lambda_e$.

Since a potential D2D UE may choose either cellular or D2D mode, the spectral efficiency of an arbitrary potential D2D UE is the average of the two operating modes. By using Theorem 1 and Lemmas 2, 3, the average term $E \left[ \log \left( 1 + \frac{G_0}{I + N} \right) \right]$ in (22) can be calculated and the spectral...
efficiency of the D2D and cellular modes can be evaluated as follows.

**Theorem 2:** For an overlaid D2D network, the spectral efficiency of a cellular UE (Rc), a potential D2D UE operating in D2D modes (Rd), and a potential D2D UE (\(\tilde{R}_d\)) are given by

\[
R_c = \frac{(1 - \beta)}{\lambda_c} \left(1 - \mathbb{E} \left[ \log \left(1 + \frac{G_0\xi}{I_c + N}\right) \right] \right),
\]

\[
R_d = \frac{\beta c}{\lambda_d} \left(1 - \mathbb{E} \left[ \log \left(1 + \frac{G_0\xi}{I_d + N}\right) \right] \right),
\]

\[
\tilde{R}_d = P(L_d > \theta) R_c + P(L_d \leq \theta) R_d,
\]

where \(\beta\) is the spectrum partition factor, \(c\) is the ALOHA transmit probability, \(\xi\) is a fitting parameter for the D2D link length distribution and \(P(L_d > \theta) = 1 - \exp(-\beta L_d^{\theta - 1})\) from (6). The average term \(\mathbb{E} \left[ \log \left(1 + \frac{G_0\xi}{I_c + N}\right) \right]\) can be evaluated using Theorem 1 as follows

\[
\mathbb{E} \left[ \log (1 + \gamma) \right] = \left\{ \begin{array}{ll}
\frac{1}{\mu^\kappa} \sum_{n=0}^{\infty} \frac{\mu^k n! \phi_{\kappa;\mu}(n)}{n!} & \text{for } \kappa - \mu \\
\sum_{n=0}^{\infty} \alpha_n n! \phi_{\eta;\mu}(n) & \text{for } \kappa - \mu
\end{array} \right.
\]

(24)

where \(\alpha_n\) is defined in (21), \(\phi_{\kappa;\mu}(n)\) and \(\phi_{\eta;\mu}(n)\) are defined in Theorem 1, \(\lambda_k(s)\) is given by (12) for the D2D mode and (15) for the cellular mode. The derivative terms \(g_i(s)\) for the logarithm function \(g(x) = \log(1 + \text{SINR})\) can be evaluated as follows

\[g_i(s) = \frac{1}{\mu^k (1+s)} \left( \frac{1}{\mu^k} - \frac{1}{(1+s)^k} \right)\]

(25)

**Remark 3:** We observe that \(R_c\) in (23) is a decreasing function of the spectrum partition factor \(\beta\), whereas \(R_d\) and \(\tilde{R}_d\) are increasing functions of \(\beta\). For a UE operating in cellular mode, \(R_c\) is an increasing function of the mode selection threshold \(\theta\); Given a large \(\theta\), more UEs will choose the D2D mode and the average number of the cellular UEs in a cell \(\mathbb{E}[N]\) will decrease. On the other hand, \(R_d\) is a decreasing function of \(\theta\) due to the increased D2D interference. Since \(\tilde{R}_d\) is the average of \(R_c\) and \(R_d\), \(\tilde{R}_d\) is concave function of \(\theta\). (See Section VI).

**Remark 4:** Theorem 2 can be applied to the general case when different types of fading affect the intended and interfering links. For example, if the fading observed in the intended link is \(\kappa\)-\(\mu\) distributed and that of the interference link is \(\eta\)-\(\mu\) distributed, then the average of the logarithm function can be evaluated using (24) where the Laplace transform of the interference \(\lambda_k(s)\) is given by (12) and (13) for the D2D mode (or (15) and (17) for the cellular mode).

**Remark 5:** Theorem 1 can be utilized to evaluate any performance measures that are represented as a function of SINR. The analytic function \(g(x)\) and the corresponding \(g_i(s)\) for several performance measures are summarized in Table III, where one can substitute \(i = \mu + n\) to \(g_i(x)\) for the \(\kappa\)-\(\mu\) distribution and \(i = 2\mu + 2n\) for the \(\eta\)-\(\mu\) distribution.

**TABLE III**

<table>
<thead>
<tr>
<th>Measure</th>
<th>(g(x))</th>
<th>(g_i(s))</th>
</tr>
</thead>
<tbody>
<tr>
<td>Spectral efficiency</td>
<td>(\log(1 + x))</td>
<td>(\frac{1}{\mu^k} \left( \frac{1}{\mu^k} - \frac{1}{(1+x)^k} \right))</td>
</tr>
<tr>
<td>Error bound</td>
<td>(e^{-x})</td>
<td>(-i \int_0^{1} (1 - 2x) e^{-x} dx)</td>
</tr>
<tr>
<td>Higher order moments</td>
<td>(x^r)</td>
<td>(\frac{1}{\mu^k} (1+x)^r - 1)</td>
</tr>
</tbody>
</table>

(26)

**B. Outage Probability**

The outage probability is defined for the D2D and cellular mode as follows

\[P_o(T_o) = \begin{cases} 
\mathbb{P} \left( \frac{G_0}{\lambda_d + N} < T_o \right) & \text{for D2D mode,} \\
\mathbb{P} \left( \frac{G_0}{\lambda_c + N} < T_o \right) & \text{for Cellular mode,}
\end{cases}\]

(27)

with a predefined SINR threshold \(T_o\). Although Theorem 1 presents a generalized framework to calculate any performance measure that is represented as an analytic function of the SINR, it can not be used for evaluating SINR distribution based performance measures, such as the outage probability or rate coverage probability. For the outage probability, \(g(x) = \mathbb{I}(x < T_o)\) is a step function and its higher order derivative is an unbounded impulse signal.

Instead of using Theorem 1, we use the series representation of the \(\kappa\)-\(\mu\) and \(\eta\)-\(\mu\) fading distributions and employ Campbell’s theorem [12], [14] to represent the SINR distribution in terms of the Laplace transform of the aggregate interference as follows. The SINR distributions \(\mathbb{P}(\text{SINR} < T_o)\) in (25) can be evaluated as

\[
\sum_{n=0}^{\infty} \sum_{m=0}^{\infty} b_{n,m}^\kappa e^{-T_o(\xi+\eta)^n+\mu} \left( \frac{T_o(\xi+\eta)^n+\mu}{\Omega \mu} \right)^{n+m+1} e^{-\frac{T_o(\xi+\eta)^n+\mu}{\Omega \mu}} \]

(28)

for \(\kappa\)-\(\mu\) fading, where we applied \(b_{n,m}^\kappa \mu = \frac{n!}{(n+m)!} e^{-\mu} \) and \(b_{n,m}^\eta \mu = \frac{n!}{(n+m)!} e^{-\mu} \). The term \(\mathbb{E}_t \left[ e^{-\frac{T_o(\xi+\eta)^n+\mu}{\Omega \mu}} \right]\) in (26) can be evaluated as follows

\[
\mathbb{E}_t \left[ e^{-\frac{T_o(\xi+\eta)^n+\mu}{\Omega \mu}} \right] = e^{-\frac{T_o(\xi+\eta)^n+\mu}{\Omega \mu}} L_{\xi+\eta} \left( \frac{T_o}{\Omega} \right),
\]

(29)

where \(L_j (x)\) is derived in Lemma 2 for the D2D link and Lemma 3 for the cellular link. Therefore, the outage probability of an overlaid D2D network is derived as follows

\[\mathbb{P}(\text{SINR} < T_o) = \sum_{n=0}^{\infty} \sum_{m=0}^{\infty} b_{n,m}^\kappa e^{-T_o(\xi+\eta)^n+\mu} (T_o) \quad \text{for } \kappa - \mu,\]

(30)

\[\sum_{n=0}^{\infty} \sum_{m=0}^{\infty} b_{n,m}^\eta e^{-T_o(\xi+\eta)^n+\mu} (T_o) \quad \text{for } \eta - \mu,\]

(31)

7The detailed proof of Table III is given in [42].
Fig. 2. Spectral efficiency of an overlaid D2D network for various channel parameters ($\kappa, \mu, \tilde{w}$); (a)-(h) assume $\lambda_b = \frac{1}{\pi 500^2}$, $\lambda = 10\lambda_b$, $\epsilon = 0.8$, $\alpha = 4$, $\beta = 0.2$, $\theta = 100m$, $q = 0.2$ and $N = \frac{N_0}{\tau} = -60$ (dBm/Hz).
where (61) is used for the $\eta$-$\mu$ fading and $\chi^{(t)}(T_o)$ denotes the higher order derivative terms as

$$
\chi^{(t)}(T_o) = (-1)^{t} \frac{d^{t}}{d\Omega^{t}} \left[ e^{-\frac{N \Omega}{\bar{\lambda} b}} L_t \left( \frac{s T_o}{\Omega} \right) \right] |_{s=1}.
$$

(29)

The outage probability of the D2D mode (or cellular mode) can be evaluated by substituting $L_t(s)$ from Lemma 2 (or $L_t(s)$ from Lemma 3) into (29). The higher order derivative in (29) can be numerically evaluated by using Faa di Bruno’s formula [45], as used in some related studies [31], [46].

VI. NUMERICAL RESULTS

In the following, we compare numerical results for different fading models. All of the simulations were carried out using MATLAB with the following parameters: BS node intensity $\lambda_b = \frac{1}{\pi 300}$, UE node intensity $\lambda = \frac{10}{\pi 300}$, ALOHA transmit probability $\epsilon = 0.8$, path-loss exponent $\alpha = 4$, spectrum partition factor $\beta = 0.2$, mode selection threshold $\theta = 100 m$, probability of potential D2D UEs $q = 0.2$, and effective noise $N = \frac{N_0}{\bar{\lambda} b} = -60$ (dBm/Hz), where $N_0$ is the noise spectral density and $\tau$ is the reference path-loss at a unit distance. Without loss of generality, we assume identical fading parameters across the intended and interference links.

We have assumed $\bar{\xi} = \bar{\lambda}$ in Figs 2(a)-(d) and respectively used $\bar{w} = 1$ for Fig. 2(a), $\mu = 1.2$ for Fig. 2(b), $\mu = 1$ for Fig. 2(c) and $\kappa = 1$ for Fig. 2(d). We observe that a dominant LOS component (large $\kappa$), a large number of scattering clusters (large $\mu$) and higher average of the channel coefficients (large $\bar{w}$) collectively achieve a higher spectral efficiency. In a weak LOS condition, the D2D link achieves higher spectral efficiency than the cellular links because, on average, D2D links have a closer transmission range than cellular links. In a strong LOS condition, the cellular links achieve higher spectral efficiency than the D2D links. In this case, the rate performance of the D2D link deteriorates due to the increased interference power from closely located D2D UEs. On the other hand, cellular links employ orthogonal medium access to ensure only one active transmitter within the cell at a given resource block. The received signal of the cellular links is protected against the elevated interference power, achieving higher cellular rate than the D2D links. We also note that the spectral efficiency is an increasing function of $\bar{w}$. Here, the rate increment of the D2D link is notable over the whole range of $\kappa$, whereas the increment of cellular link is distinguishable only after $\kappa \geq 5$.

In Figs 2(e)-(h), we have assumed $\bar{\xi} = 1$ and $\bar{w} = 1$ and used $\mu = 1$ for Fig. 2(e), $\kappa = 2$ for Fig. 2(f), $\mu = 1$ for Fig. 2(g) and $\kappa = 1$ for Fig. 2(h), respectively. Since the spectral efficiency of a cellular UE $R_c$ is a decreasing function of the spectrum partition factor $\beta$ (and $\bar{R}_d$ is an increasing function of $\beta$), there is a crossover point $\beta^*$ between the spectral efficiencies of the D2D and cellular link, which depends on the fading parameters. On average, D2D links have a closer transmission range than cellular links, hence if a minimum amount of spectrum is allocated to the D2D link, which is $\beta \geq \beta^*$, D2D UEs achieve higher transmission rates than cellular UEs. On the other hand, if $\beta < \beta^*$, D2D transmission does not have enough radio resources to achieve rate gains against the cellular link. As the number of scattering clusters increases, the spectral efficiency of the cellular UE becomes larger than that of the D2D UE and the crossover point $\beta^*$ shifts towards the right.

Figs 2(g)-(h) show the effect of the mode selection threshold $\theta$ on the spectral efficiency for $\kappa$-$\mu$ fading. As shown in the figure, increasing $\theta$ results in less potential D2D UEs choosing to operate in the cellular mode, leading to a higher average rate for the cellular link. The spectral efficiency of a potential D2D UE in D2D mode $R_d$ is a decreasing function of $\theta$ due to the increased co-channel interference over the D2D link. Since the spectral efficiency of a potential D2D UE $R_d$ is a weighted average of $R_c$ and $R_d$, $\bar{R}_d$ is concave function of $\theta$ as indicated in Figs 2(g)-(h).

VII. CONCLUSION

In this paper, we have considered a D2D network overlaid on an uplink cellular network, where the locations of the mobile UEs as well as the BSs are modeled as PPP. In particular, we exploited a novel stochastic geometric approach for evaluating the D2D network performance under the assumption of generalized fading conditions described by the $\kappa$-$\mu$ and $\eta$-$\mu$ fading models. Using these methods, we evaluated the spectral efficiency and outage probability of the overlaid D2D network. Specifically, we observed that the D2D link provides higher rates than those of the cellular link when the spectrum partition factor was appropriately chosen. Under these circumstances, setting a large mode selection threshold will encourage more UEs to use the D2D mode, which increases the average rate at the cost of a higher level of interference and degraded outage probability. However, for smaller values of the spectrum partition factor, the D2D link has smaller rates than those of the cellular link. In terms of the fading parameters, a dominant LOS component (large $\kappa$) or a large number of scattering clusters (large $\mu$) improve the network performance, i.e., a higher rate and lower outage probability are achieved. Finally, we also provided numerical results to demonstrate the performance gains of overlaid D2D networks compared to traditional cellular networks, where the latter corresponds to the $\beta = 0$ case.

APPENDIX I

For conciseness, in this appendix, we summarize the operational equalities of some special functions, which are used in this paper.\textsuperscript{8} The following properties hold for non-negative real constants $x$ and $y$

$$
\Gamma(x)\Gamma\left(x + \frac{1}{2}\right) = 2^{1-2x}\sqrt{\pi}\Gamma(2x),
$$

$$
\frac{\Gamma\left(x + \frac{1}{2}\right)}{\Gamma(y + 1)\Gamma(x - y + 1)} = \frac{\Gamma\left(x + 1\right)}{\Gamma\left(y + 1\right)\Gamma\left(x - y + 1\right)}.
$$

(27)

$$
\Gamma(1+x)\Gamma(1-x) = \frac{1}{\text{sinc}(x)}, \quad \Gamma\left(x + \frac{1}{2}\right) = \sqrt{\pi}.
$$

(28)

\textsuperscript{8}Most of the expressions in Appendix I were introduced in [47], except for (38) and (39), which were proved in [48].
The following properties of the hypergeometric function hold for real constants \( a, b \) and \( c \):

\[
1_F_1(a; b; t) = e^t 1_F_1(b - a; b; -t),
\]

\[
2_F_1(a; b; c; z) = (1 - z)^{-a} 2_F_1 \left( a, c - b; c; \frac{z}{z - 1} \right),
\]

\[
\int_0^\infty t^{a-1} e^{-t} 1_F_1(a; b; -t) dt = \Gamma(a) e^{-a} 2_F_1 \left( a, a; b; -\frac{1}{c} \right),
\]

\[
((a - b)z + c - 2a) 2_F_1(a; b; c; z) = (c - a) 2_F_1(a - 1, b; c; z) + a(z - 1) 2_F_1(a + 1; b; c; z),
\]

\[
\int_0^\infty x^{d-1} e^{-bx} \gamma(c, ax) dx = \frac{a^e \Gamma(d + c)}{c(a + b)^d + c} 2_F_1 \left( 1, d + c; c + 1; \frac{a}{a + b} \right),
\]

where (32) holds for \( a > 0 \) and \( c > 0 \), (34) holds for \( a + b > 0 \), \( b > 0 \), and \( c + d > 0 \). The modified Bessel function \( I_v(x) \) and incomplete gamma function \( \gamma(s, x) = \int_0^x t^{s-1} e^{-t} dt \) can be represented by the hypergeometric function with arbitrary positive real constants \( v, s, b \) as follows:

\[
I_v(2\sqrt{bt}) = \frac{0_F_1(v; b; t)}{\Gamma(v)} (bt)^{\frac{v}{2}},
\]

\[
\gamma(s, x) = s^{-1} x^{e-\frac{x}{s}} 2_F_1(1; 1 + s; x),
\]

where

\[
I_v(x) = \sum_{n=0}^{\infty} \frac{1}{n! \Gamma(n + v + 1)} \left( \frac{x}{2} \right)^{2n+v},
\]

\[
\gamma(s, x) = \sum_{n=0}^{\infty} \frac{x^{s+n} e^{-x}}{\Gamma(s + n + 1)},
\]

\[
0_F_1(v; a; b) = \lim_{v \to 0} \Gamma(v) 2_F_1 \left( a; v; \frac{b}{a} \right).
\]

\[F_2(a; \beta; \gamma; y, x, y) = \sum_{m=0}^{\infty} \sum_{n=0}^{\infty} \frac{(a)_{m+n} (\beta)_{m} (\gamma)_{n}}{m! n!} y^{m} x^{n}.\]

Appell’s function can be reduced to the hypergeometric function using the following properties:

\[F_2(a; \alpha, \beta; c; 0, 0) = 2_F_1(d; a; c; 0, 0), \quad F_2(a; \alpha, \beta; c; 0, y) = 2_F_1(d; a; c; y, 0), \]

\[\int_0^\infty t^{d-1} e^{-ht} 1_F_1(a; b; -t) dt = \frac{h^{-d} \Gamma(d) 2_F_1(d; a, b; b; k')}{h},
\]

Gaussian-Laguerre quadratures can be used to evaluate the integral for a given analytic function \( g(x) \) as

\[\int_0^\infty e^{-x} g(x) dx = \sum_{n=1}^{N} w_n f(x_n) + R_N,\]

where \( x_n \) and \( w_n \) are the \( n \)-th abscissa and weight of the \( N \)-th order Laguerre polynomial, respectively. The remainder \( R_N \) rapidly converges to zero.

**Appendix II**

In this appendix, we provide a proof for Lemma 1. First, the transmit power of a cellular UE is \( P_c = L_c^{a} \) and the pdf of \( L_c \) is given by (4). Then, the \( i \)-th moment of \( P_c \) is given by

\[E \left[ P_c^i \right] = \int_0^\infty (x^a)^i f_L(x) dx = 2\pi \lambda b \int_0^\infty x^{a+1} e^{-\lambda x^2} dx \]

\[= (\lambda b)^{-\frac{1}{2}} \Gamma \left( \frac{1}{2} + 1 \right), \]

where we applied a change of variable, i.e., \( \lambda x^2 = t \), and used the notion of the gamma function \( \Gamma(t) = \int_0^\infty x^{t-1} e^{-x} dx \) in the last equality. Similarly, the transmit power of D2D mode is \( P_d = L_d^{a} \), given that the criterion \( L_d \leq \theta \) is met. Then, the \( i \)-th moment of \( P_d \) is given by

\[E \left[ P_d^i \right] = E \left[ L_d^{a} \right] \left( \int_0^\theta x^{a+1} f_L(x) P(L_d \leq \theta) dx \right) \]

\[= \frac{2\pi \lambda'}{1 - e^{-\lambda' \theta^2}} \left[ \frac{\gamma \left( \frac{1}{2} + 1, \lambda' \theta^2 \right)}{1 - e^{-\lambda' \theta^2}} \right], \]

where we substituted (6) in the second equality, applied a change of variable, i.e., \( \lambda' x^2 = t \), and used the notion of the lower incomplete gamma function \( \gamma(s, x) = \int_0^x t^{s-1} e^{-t} dt \) with \( \lambda' = \frac{x}{\lambda} \) in the last equality. Since a potential D2D UE may choose either cellular or D2D mode, the average transmit power of an arbitrary potential D2D UE is the average of the two operating modes as follows

\[E \left[ P_d' \right] = P(L_d \leq \theta) E \left[ P_d^i \right] + P(L_d > \theta) E \left[ P_c^i \right]. \]

which we obtain (7) by substituting (42) and (43) in (44). This completes the proof.

**Appendix III**

In this appendix, we provide a proof for Lemma 2. The Laplace transform of the interference at the D2D receiver \( L_d(s) \) is evaluated as follows

\[\exp \left( -2\pi \lambda_d \int_0^\infty \left( 1 - E \left[ sG \right] \right) r^{a-1} \right) dr, \]

where we considered the location independent thinning by ALOHA transmit probability \( \varepsilon \) and used the probability generating functional (PGFL) of PPP [12]. Then, by applying a change of variable, i.e., \( sG = r^{a-1} \), and integration by parts, the Laplace transform \( L_d(s) \) simplifies as

\[\exp \left( -2\pi \varepsilon \lambda_d \left[ sG \right] \right) \int_0^\infty \delta r^{a-1} (1 - e^{-r}) dr \]

\[= \exp \left( -2\pi \varepsilon \lambda_d \left[ sG \right] \right) \left[ \int_0^\infty \delta r^{a-1} (1 - e^{-r}) dr \right] \]

\[= \exp \left( -2\pi \varepsilon \lambda_d \left[ sG \right] \right) \left[ \int_0^\infty \delta r^{a-1} (1 - e^{-r}) dr \right]. \]
where we used the gamma function $\Gamma(t) = \int_0^\infty x^{t-1}e^{-x}dx$ in the last equality. Hence, (13) for $\kappa-\mu$ fading can be obtained by substituting $\lambda_d = \pi P(L_d \leq \theta)\lambda_s$, (7), (9) into (46) as follows

$$c_d = \pi \epsilon_d \lambda_d \Gamma(1 - \delta) \mathbb{E}\left[L_d^\delta\right] \mathbb{E}\left[G^\delta\right]$$

$$= \frac{q \epsilon \delta}{\text{sinc}(\delta)} \cdot \gamma \left(2, \lambda \pi \theta^2 \cdot \xi^{-1}\right) F_1(\mu + \delta; \mu; \mu \kappa) \times \left(\frac{d}{(1 + \kappa)}\right)^\delta \left(\mu + \delta - 1\right).$$

(47)

where we used (30) in the last equality. Similarly, the Laplace transform of $I_3$ for the $\eta-\mu$ distribution can be evaluated by using (10). This completes the proof.

APPENDIX IV

In this appendix, we provide a proof for Lemma 3. The Laplace transform of the interference at the cellular BS $L_i(s)$ is evaluated as follows [13], [38]

$$L_i(s) = \exp(-2\pi \lambda_b \phi(s)),$$

(48)

where $\phi(s)$ is defined as follows

$$\phi(s) = \int_0^\infty \left(1 - e^{-\lambda_b x^2}\right) \mathbb{E}\left[1 - e^{-sGt^\nu x}a\right] dx$$

$$= \int_0^\infty \int_0^x r e^{-\lambda_b r^2} \mathbb{E}\left[1 - e^{-sGr^\nu x}a\right] dr \cdot dx$$

$$= \int_0^\infty r e^{-\lambda_b r^2} \int_{r}^\infty x \cdot \mathbb{E}\left[1 - e^{-sGr^\nu x}a\right] dx dr$$

$$= I_1 \cdot I_2,$$

(49)

where the PGFL of non-homogeneous PPP with intensity function $\lambda_b \left(1 - \exp\left(-\pi \lambda_b d^2\right)\right)$ is used in the first equality, (5) is applied in the second equality. Fabini’s theorem [49] is utilized to change the order of integration in the third equality and a change of variable, i.e., $(r/x)^2 = t$, is used in the last equality. $I_1$ and $I_2$ in (49) represent the following integrals

$$I_1 = \int_0^\infty r^3 \exp\left(-\pi \lambda_b r^2\right) dr,$$

$$I_2 = \mathbb{E}_G\left\{\int_0^1 t^{-2} \left(1 - \exp\left(-sGr^\nu t\right)\right) dt\right\},$$

(50)

which convert the double integral into a multiplication of two single integrals that are independent to each other.

The first integral $I_1$ in (49) can be evaluated by using a change of variable, i.e., $\pi \lambda_b r^2 = t$, and the definition of the gamma function $\Gamma(t) = \int_0^\infty x^{t-1}e^{-x}dx$ as follows

$$I_1 = \frac{1}{2(2\pi \lambda_b)^2} \int_0^\infty t \exp(-t) dt = \frac{\Gamma(2)}{2(2\pi \lambda_b)^2} = \frac{1}{2(2\pi \lambda_b)^2}. $$

(51)

The second integral $I_2$ can be simplified by using a change of variable, i.e., $sGr^\nu = u$, integration by parts and the definition of the lower incomplete gamma function $\gamma(s, x)$ as follows

$$I_2 = \mathbb{E}_G\left\{(sG)^\delta \int_0^{sG} \delta u^{-\delta - 1} \left(1 - \exp\left(-u\right)\right) du\right\}$$

$$= \mathbb{E}_G\left\{(sG)^\delta \gamma(1 - \delta, sG) - (1 - \exp\left(-sG\right))\right\}$$

$$= I_3 + L_G(s) - 1,$$

(52)

where $L_G(s)$ represents the Laplace transform of the channel coefficient $G$ and $I_3 \triangleq \mathbb{E}_G\left[sG\right]^\delta \gamma(1 - \delta, sG)$.

For $\kappa-\mu$ fading, the integral $I_3$ can be evaluated as follows

$$I_3 = \left(\frac{\pi}{\mu(1 + \kappa)}\right)^\delta e^{-\mu k} \times \int_0^\infty t^{\delta + 2n} \gamma(1 - \delta, s\Omega) t^{\mu - 1} (2\sqrt{\mu kt}) e^{-t} dt$$

$$= \frac{s \Omega e^{-\mu k}}{(1 - \delta) \Gamma(\mu)} \times \int_0^\infty t^{\mu - 1} \gamma(1, s\Omega) F_1(1; 2 - \delta; s\Omega) dt$$

$$= \frac{s \Omega e^{-\mu k}}{(1 - \delta) \Gamma(\mu)} \times \int_0^\infty \left(\frac{\pi}{\mu(1 + \kappa)}\right)^\delta \left(\mu + \delta - 1\right).$$

(53)

where we applied (9) with $\Omega_{\kappa, \mu} = \frac{\pi}{\mu(1 + \kappa)}$ in the first equality, utilized (35) in the second equality, then used (37), (39) and (40) in the last equality. By substituting $L_G(s)$ of (9), (51), (52), (53) to (48), we obtain (16) for the $\kappa-\mu$ distribution. Similarly for $\eta-\mu$ fading,

$$I_3 = \frac{s \Omega}{\mu(1 - \delta) \Gamma(\mu)} \sum_{n=0}^{\infty} \left(\frac{H}{h}\right)^{2n} \left(\mu + 1\right)^n \Omega(2n + 2\mu)$$

$$\times \gamma(1 - \delta, s\Omega) \sum_{n=0}^{\infty} \left(\frac{H}{h}\right)^{2n} \left(\mu + 1\right)^n$$

$$= \frac{2s \Omega H^\mu}{(1 - \delta) \Gamma(\mu)} \sum_{n=0}^{\infty} \left(\frac{H}{h}\right)^{2n} \left(\mu + 1\right)^n$$

$$\times \gamma(1 - \delta, s\Omega)$$

where we used the series representation of the Bessel function (36) in the first equality, applied (34) in the second equality and used a change of variable, i.e., $n + \mu = n'$, in the last equality. The corresponding Laplace transform of interference over $\eta-\mu$ fading can be obtained by substituting $L_G(s)$ of (10), (51), (52), (54) to (48). This completes the proof.

APPENDIX V

In this appendix, we provide a proof for Theorem 1. First, we consider $\kappa-\mu$ fading and obtain the series representation of the $\kappa-\mu$ distribution by using (9) and (36) as follows

$$f_G(x) = \frac{1}{e^{\mu k}} \sum_{n=0}^{\infty} \frac{(\mu k)^n}{n!} \frac{\Omega^{-n}\mu}{\Gamma(n + \mu)} x^{n-1} \exp\left(-\frac{x}{\Omega}\right).$$

(55)

Then, the average of an arbitrary function of the SINR $\gamma = \frac{G_0}{I + N}$ for a given interference $I$ is

$$\mathbb{E}\left[\log\left(\frac{G_0}{I + N}\right)\right] = \frac{1}{e^{\mu k}} \sum_{n=0}^{\infty} \frac{(\mu k)^n}{n!} \frac{\Omega^{-n}\mu}{\Gamma(n + \mu + 1)} \exp\left(-\frac{x}{\Omega}\right) \log\left(\frac{G_0}{I + N}\right)$$

$$= \frac{1}{e^{\mu k}} \sum_{n=0}^{\infty} \frac{(\mu k)^n}{n!} \int_0^\infty \frac{e^{-\mu x - 1}}{\Gamma(n + \mu + 1)} g(z) b^{\mu + 1} e^{-bz} dz. $$

(56)
where we applied (55) in the first equality and used a change of variable, i.e., \( \frac{x}{\mu + n} = z \) and \( b = \frac{(t+\eta)}{\Omega} \), in the second equality. The integral in (56) can be evaluated as follows

\[
\int_0^\infty z^{\mu+n-1} \frac{g(z)}{\Gamma(\mu + n)} b^{\mu+n} e^{-bz} \, dz = - \sum_{i=0}^{\mu+n-1} g_i(z) b^{i+n+1} \int_0^\infty e^{-bz} \, dz,
\]

where we applied integration by parts \( \eta \) times, defined \( g_i(z) \) in (19), and

\[
g_i(0) = \begin{cases} \frac{1}{\mu^{i+n}}, & \text{for } i < \mu + n - 1 \\frac{1}{b \cdot \Gamma(\mu + n)}, & \text{for } i = \mu + n - 1. \end{cases}
\]

Then, the average of an arbitrary function of the SINR for the \( \kappa\mu \) fading is given by

\[
\mathbb{E} \left[ g \left( \frac{G_0}{I+N} \right) \right] = \mathbb{E} \left[ \mathbb{E} \left[ g \left( \frac{G_0}{I+N} \right) | I \right] \right] = \frac{1}{\mu^{\mu+n}} \sum_{n=0}^{\infty} \frac{(\mu \kappa)^n}{n!} \phi_{\kappa,n}(n),
\]

where we applied \( \frac{1}{\mu^{\mu+n}} \sum_{n=0}^{\infty} \frac{(\mu \kappa)^n}{n!} = 1 \) Similarly, (20) can be evaluated for \( \eta\mu \) fading by repeatedly applying integration by parts. This completes the proof.

APPENDIX VI

In this appendix, we express the CDF of the \( \kappa\mu \) and \( \eta\mu \) distributions using a series representation. First, we integrate the PDF of the \( \kappa\mu \) distribution in (55) as follows

\[
F_G(x) = \int_0^x f_G(t) \, dt = \sum_{n=0}^{\infty} \frac{(\mu \kappa)^n \gamma(n+\mu, x)}{\Gamma(n+\mu)} = \sum_{n=0}^{\infty} \sum_{m=0}^{\infty} \frac{b_{\mu,m}^{\kappa,n}}{\Omega^{n+\mu+m}} \exp \left( - \frac{x}{\Omega} \right),
\]

where we used the definition of the incomplete gamma function \( \gamma(s,x) = \int_0^x t^{s-1} e^{-t} \, dt \) in the second equality, applied (36) in the third equality, denoted \( \Omega_{\mu} = \frac{\mu}{\mu + (1+\kappa)} \) and \( b_{\mu,m}^{\kappa,n} \triangleq \frac{\frac{(\mu \kappa)^n \gamma(n+\mu, x)}{\Gamma(n+\mu)}}{\frac{\mu^{\mu+n}}{\mu^{\mu+n} + \mu (1+\kappa)}} \). Similarly, for the \( \eta\mu \) distribution, the series representation of its CDF is given by

\[
F_G(x) = \sum_{n=0}^{\infty} \frac{b_{\mu,m}^{\eta,n}}{\Omega^{n+\mu+m}} \exp \left( - \frac{x}{\Omega} \right),
\]

where \( b_{\mu,m}^{\eta,n} \triangleq \frac{1}{\mu^n} \frac{\gamma(n+\mu-1)}{\Gamma(n+\mu)} \frac{(\mu \kappa)^n}{\Omega^{n+\mu+m}} \).

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